

An Analytic Model for the e-Cloud in a Magnetic Field Dominated Ring



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Outline

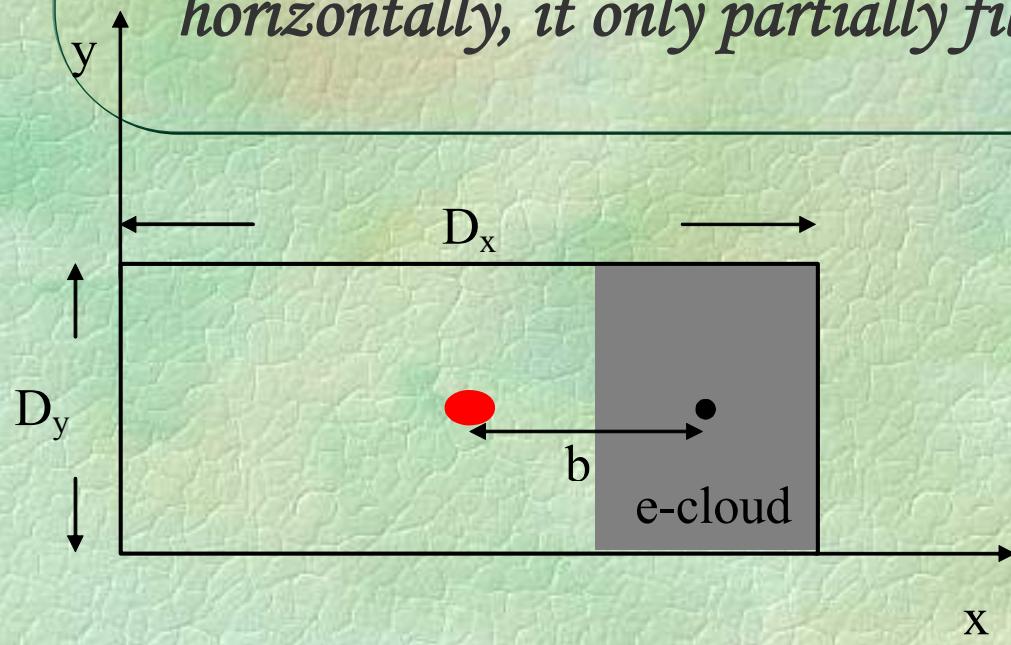
- ✓ *Model Description*
- ✓ *Vertical tune-shift*
- ✓ *Build-up*
- ✓ *Life-time*
- ✓ *Theory and experiment*
- ✓ *Summary and conclusions*

Model Description

- ✓ Along the ring each bunch experiences , in **average**, a certain cloud density:

$$q(s) = \sum_{n=-\infty}^{\infty} q_n \exp(j2\pi ns/C) \approx q_0$$

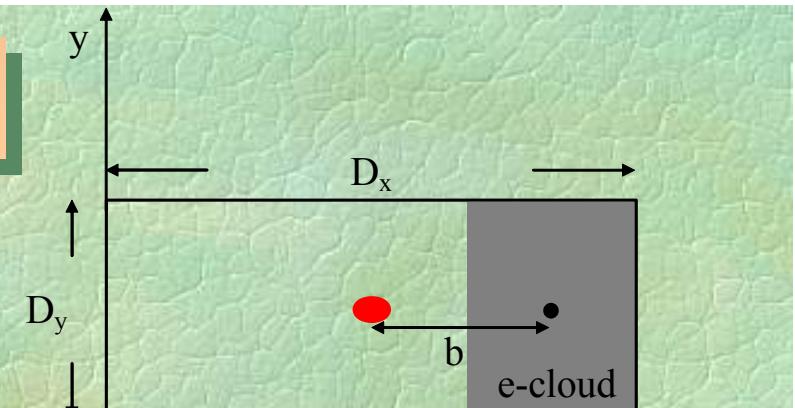
- ✓ This density **varies** from bunch to bunch.
- ✓ **Effective life-time** is the average life-time of all bunches
- ✓ The cloud is assumed to have **vertical symmetry** but horizontally, it only partially fills the chamber.



$$n_{ec} = \frac{q_0}{e} \frac{1}{D_x D_y C}$$

Circumference

Vertical Force of the E-cloud



- ✓ Focusing for positrons ($p=+e$); defocusing for electrons ($p= -e$).

$$\left[\frac{d^2}{ds^2} + \frac{1}{\beta_y^2} \right] \delta y = \frac{p E_y}{mc^2 \gamma} = \frac{\delta y}{\beta_{ec}^2}$$

$$\frac{1}{\beta_{ec}^2} = -\frac{p}{m_0 c^2 \gamma} \frac{Q_0}{4\pi \epsilon_0 C} \frac{f_E(\bar{b} = 2b/D_x)}{D_y^2}$$

- ✓ Uniform cloud. Its center-of-mass at a distance b from center

$$f_E(\bar{b}) = \frac{(32/\pi)}{1-\bar{b}} \sum_{n_x, n_y=1}^{\infty} \frac{n_y}{n_x} \frac{\cos(\pi n_x \bar{b}) - \cos(\pi n_x)}{n_x^2 (D_y/D_x) + n_x^2 (D_y/D_x)} \sin\left(\frac{\pi}{2} n_x\right) \sin^3\left(\frac{\pi}{2} n_y\right)$$

Relative vertical tune - shift

$$Q_y = \frac{1}{2\pi} \int_0^s ds' \sqrt{\frac{1}{\beta_y^2} - \frac{1}{\beta_{ec}^2}} \simeq \frac{1}{2\pi} \int_0^s ds' \frac{1}{\beta_y} - \frac{1}{2\pi} \frac{1}{2} \int_0^s ds' \frac{\beta_y}{\beta_{ec}^2}$$

$$\delta Q [\%] = 100 \times \frac{Q_y - Q_y^{(0)}}{Q_y^{(0)}} \simeq -\frac{100}{2} \frac{\oint ds \frac{\beta_y(s)}{\beta_{ec}^2(s)}}{\oint ds \frac{1}{\beta_y(s)}} \simeq -\frac{50}{\beta_{ec}^2} \frac{\oint ds \beta_y(s)}{\oint ds \frac{1}{\beta_y(s)}}$$

$$\boxed{\delta Q [\%] \simeq -50 \frac{\beta_{coupling}^2}{\beta_{ec}^2} \simeq \left\{ 50 \frac{\beta_{coupling}^2 r_e}{\gamma} \frac{D_x}{D_y} \underbrace{f_E(\bar{b})}_{EC \text{ geometry}} \right\} n_{ec}}$$

✓ The relative tune shift is proportional to the **average cloud density**, cloud geometry and the lattice **coupling parameter**

$$\beta_{coupling}^2 = \frac{\oint ds \beta_y(s)}{\oint ds \frac{1}{\beta_y(s)}} \simeq \begin{cases} 240.8 [\text{m}^2] & 12 \text{ wigglers} \\ 252.4 [\text{m}^2] & 6 \text{ wigglers} \end{cases}$$

Build-up

New-born electrons

- ✓ “Charging-up” of the chamber: **equilibrium** between generation and absorption at the wall

$$\frac{dn_{ec}}{dt} = -\frac{n_{ec}}{\tau} + \sum_{\nu=1}^N n_{\nu}^{(nb)} \delta(t - \nu T)$$

T - bunch spacing

ν – bunch index

τ – life-time

- ✓ Assumption: “new-borns” follow the initial average current of each one of the bunches

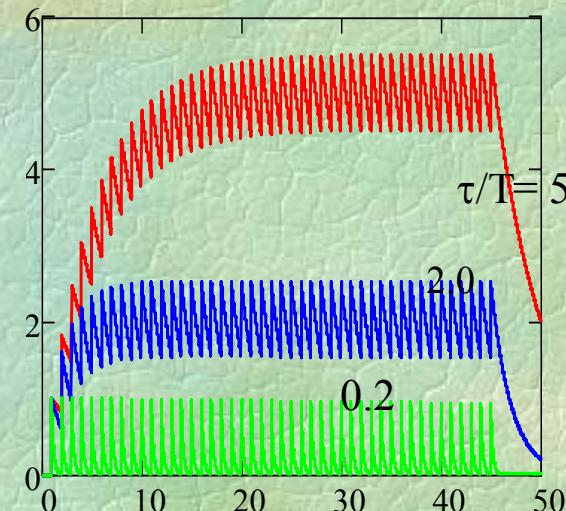
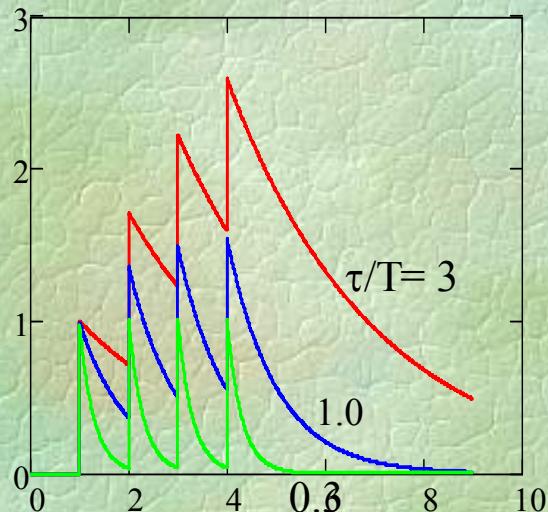
$$\frac{n_{\nu}^{(nb)}}{\langle n_{\mu}^{(nb)} \rangle_{\mu}} \simeq \frac{I_{\nu}}{\langle I \rangle}$$

- ✓ What counts is the **density experienced** by each bunch

$$n_{ec,\mu} \equiv n_{ec}(t = \mu T - 0) = n_{nb} \sum_{\nu=1}^N \frac{I_{\nu}}{\langle I \rangle} h(\mu - \nu - 0) \exp\left[-\frac{T}{\tau}(\mu - \nu)\right]$$

Average E-cloud Density

$$\frac{dn_{ec}}{dt} + \frac{n_{ec}}{\tau} = \sum_{\nu=1}^N n_{\nu}^{(nb)} \delta(t - \nu T)$$



✓ Equilibrium density determined by the *life-time* !!

$$n_{ec} \simeq \frac{\tau}{T} \left\langle n_{\nu}^{(nb)} \right\rangle$$

Life-time

$$\tau \cong \frac{D_y}{c} \sqrt{\frac{mc^2}{2E}}$$

- ✓ The life-time of a **ballistic** 100[eV] electron is 8nsec (1[eV], 80nsec)
- ✓ The dynamics of the electrons **ignoring the kick of the bunches**

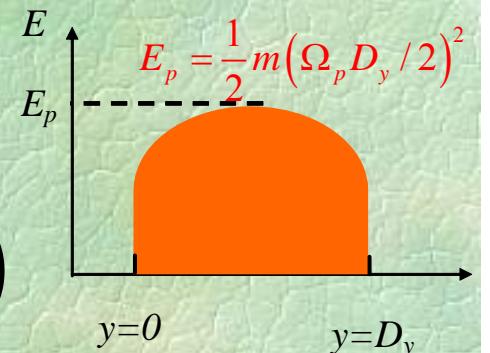
$$\left[\frac{d^2}{dt^2} - \Omega_p^2 \right] \delta y_i = 0$$

$$\delta y_i(t) = \sqrt{2E_i/m\Omega_p^2} \sinh(\Omega_p t) + (-D_y/2) \cosh(\Omega_p t)$$

$$y_i = D_y/2 + \delta y_i$$

$$\Omega_p^2 = \frac{e^2 n_{ec}}{m \epsilon_0} f_P \left(\bar{b} = 2b/D_x \right)$$

$$f_P(\bar{b}) = \frac{(2/\pi)^3}{(1-\bar{b})^2} \frac{D_x}{D_y} \sum_{n_x, n_y=1}^{\infty} \frac{n_y}{n_x^2} \frac{\sin^3(\pi n_y/2) [\cos(\pi n_x \bar{b}) - \cos(\pi n_x)]^2}{n_x^2 (D_y/D_x) + n_y^2 (D_x/D_y)}$$



Cloud as a potential barrier for the new-borns

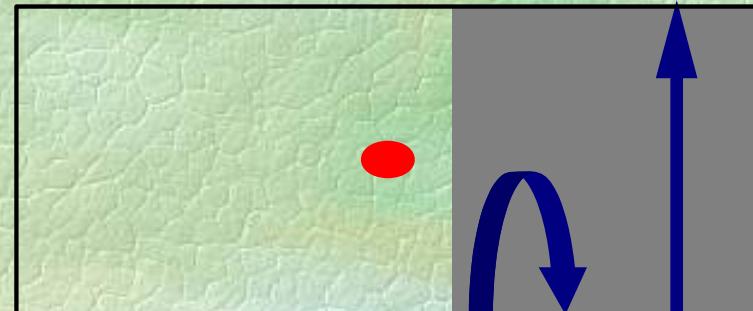
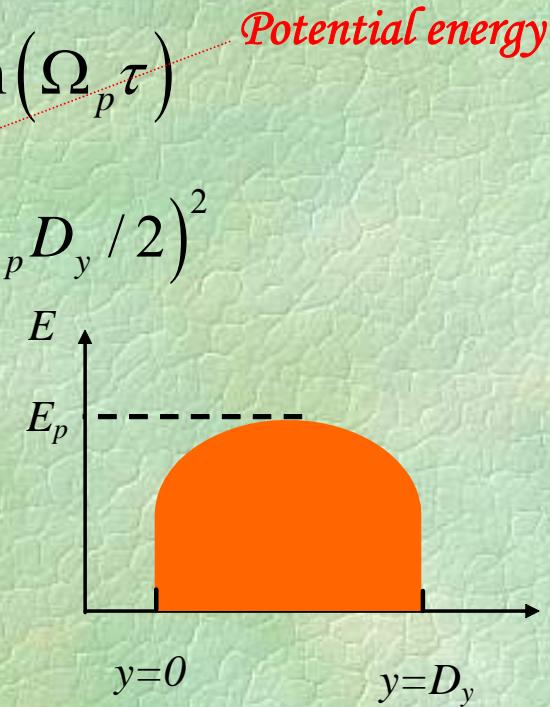
Life-time

✓ The dynamics of the electrons **ignoring the kick of the bunches**

$$D_y / 2 = \sqrt{2E_i / m\Omega_p^2} \sinh(\Omega_p \tau) + (-D_y / 2) \cosh(\Omega_p \tau)$$

$$\tau'(E_i) = \frac{2}{\Omega_p} \begin{cases} \operatorname{atanh}\left(\sqrt{E_p / E_i}\right) & E_i > E_p \\ \operatorname{atanh}\left(\sqrt{E_i / E_p}\right) & E_i < E_p \end{cases}$$

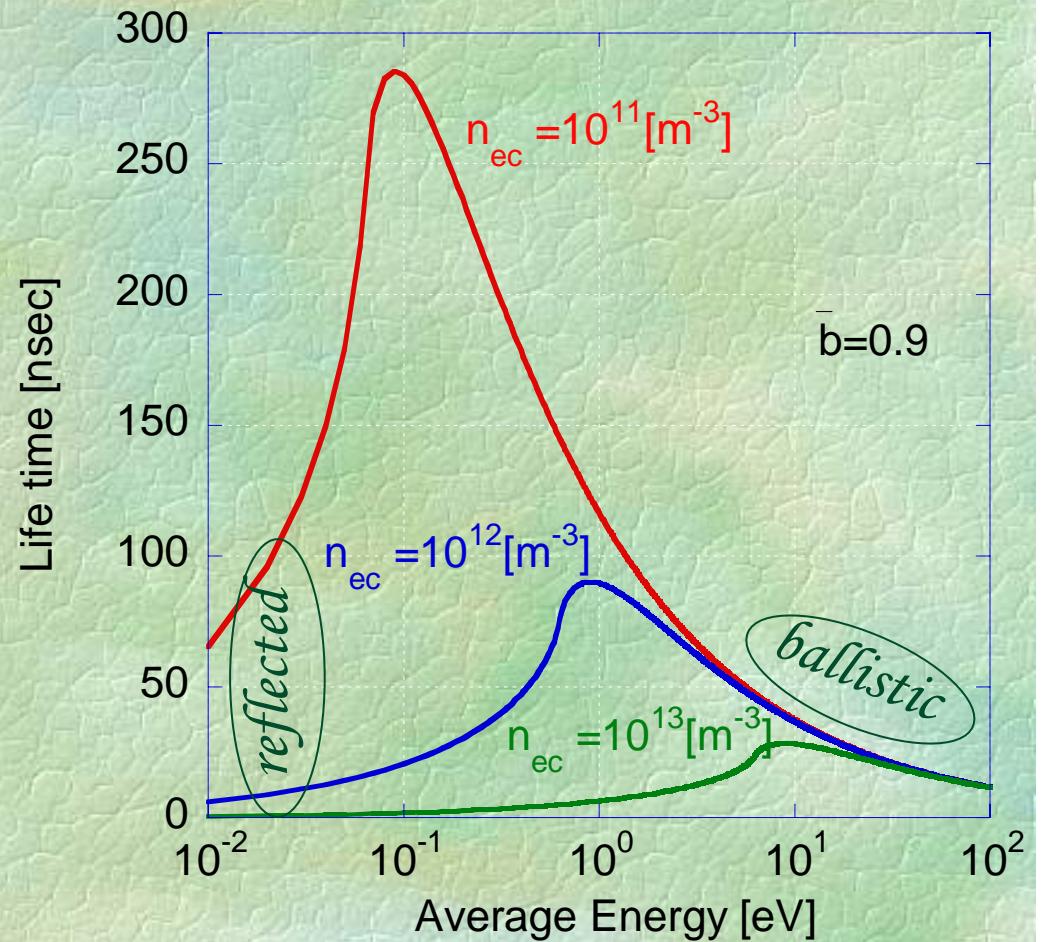
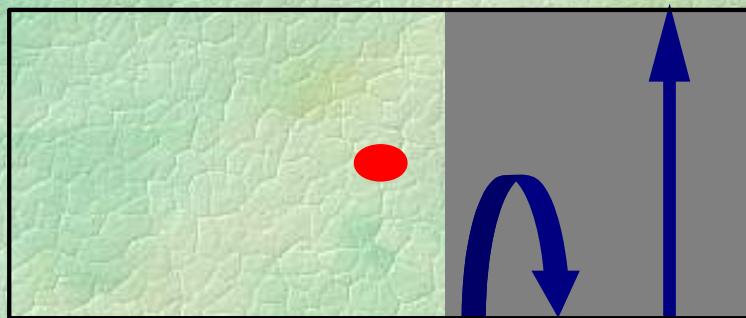
$$\boxed{\tau = \int_0^\infty dE f_{ec}(E) \tau'(E) = \frac{1}{E_0} \int_0^{E_0} dE \tau'(E)}$$



Life-time

- ✓ The dynamics of the electrons **ignoring the kick of the bunches**
- ✓ Uniform spectrum of electrons $\langle E \rangle = E_0 / 2$

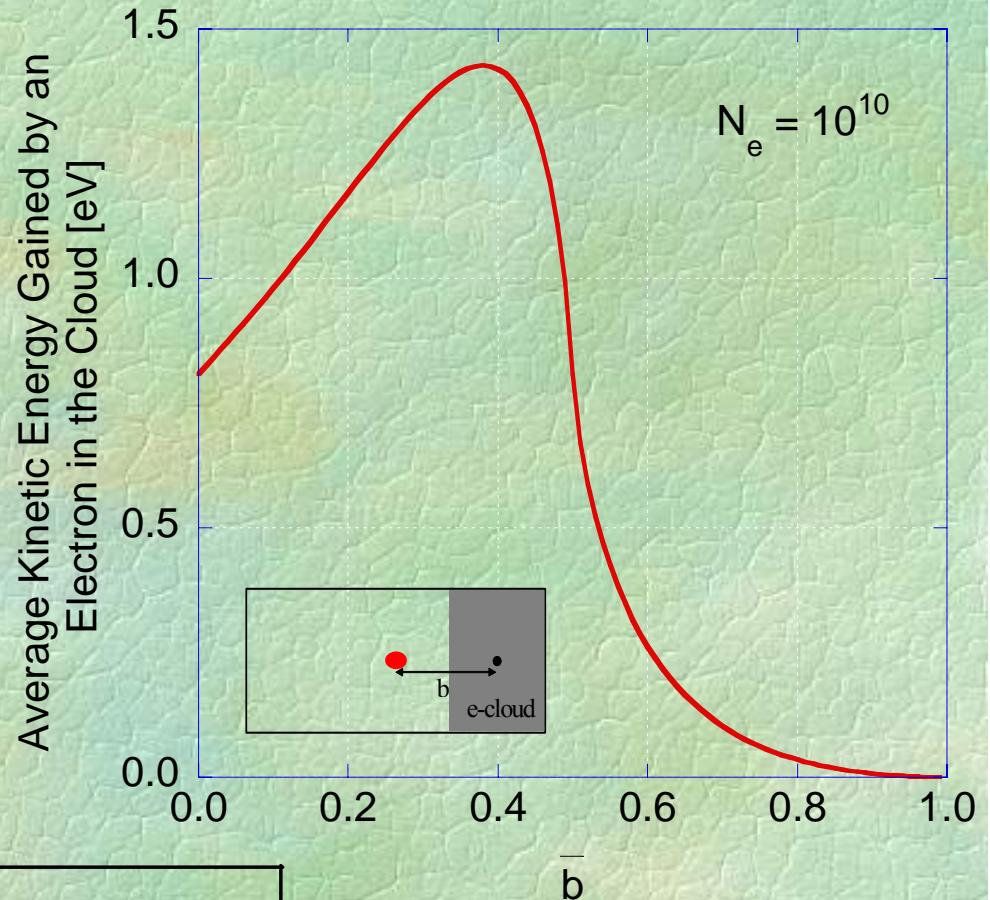
$$\tau = \frac{1}{E_0} \int_0^{E_0} dE \tau'(E)$$



Life-time

- ✓ Kick by the positrons bunch.
- ✓ Average energy transferred is less than $1.5\text{eV}!!$ No M-P.
- ✓ However, this kick may trap the electrons in case of positron bunches.

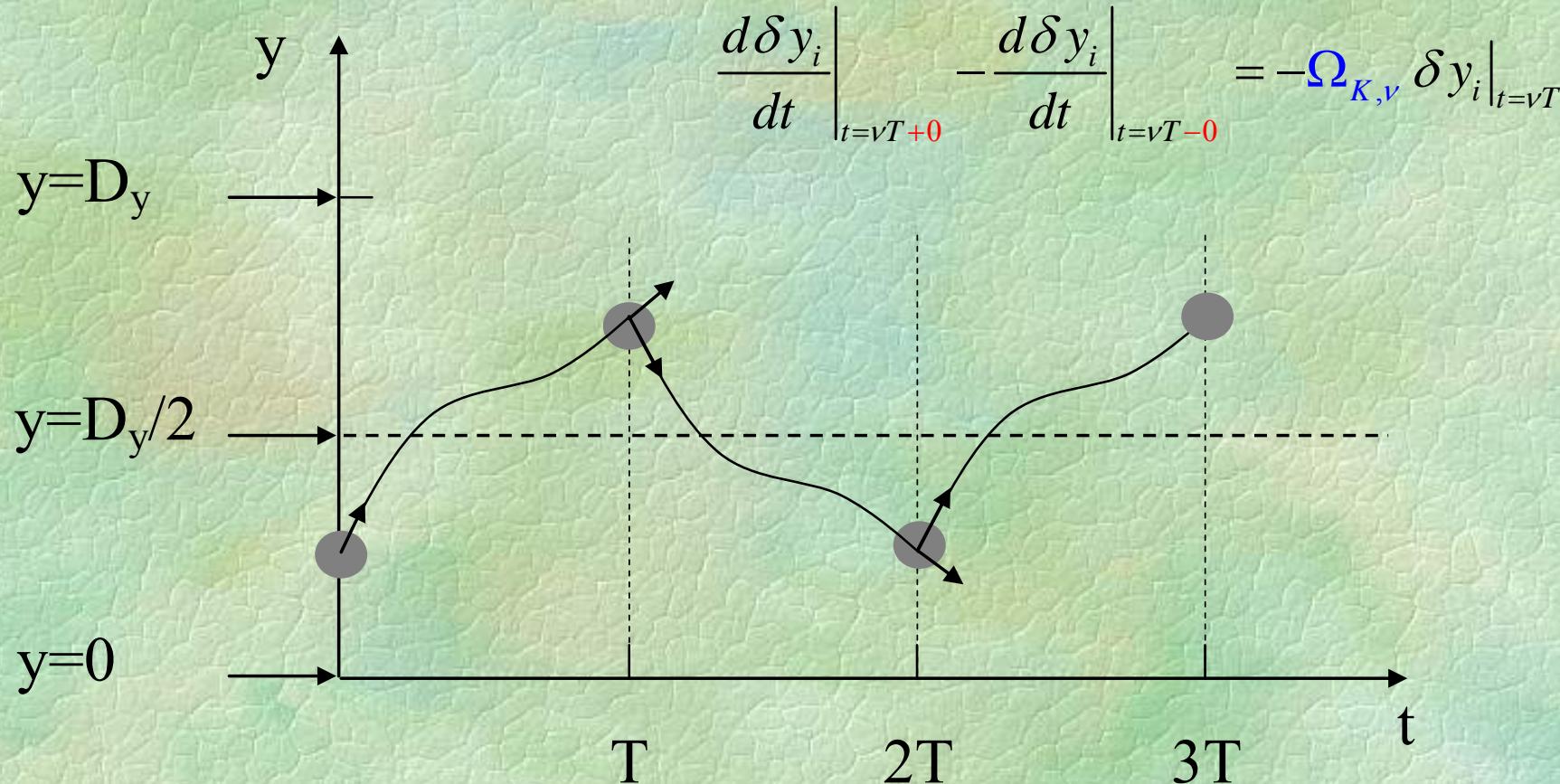
$$\left. \frac{d\delta y_i}{dt} \right|_{t=\nu T+0} - \left. \frac{d\delta y_i}{dt} \right|_{t=\nu T-0} = -\Omega_{K,\nu} \delta y_i \Big|_{t=\nu T}$$



$$\Omega_{K,\nu}^2 = \left(\frac{e^2 N_{e,\nu}}{m \epsilon_0 D_y^3} \frac{D_y}{c} \right)^2 \frac{96/\pi^4}{\left(1-\bar{b}\right)^2} \sum_{n_y=1}^{\infty} \left[\sum_{n_x=1}^{\infty} \frac{n_y}{n_x} \frac{\sin\left(\frac{\pi}{2} n_x\right) \sin\left(\frac{\pi}{2} n_y\right)}{n_x^2 \frac{D_y}{D_x} + n_y^2 \frac{D_x}{D_y}} \left[\cos\left(\pi n_x \bar{b}\right) - \cos\left(\pi n_x\right) \right] \right]^2$$

Life-time

✓ *Trapping by the positrons train: qualitative picture*



Life-time

✓ *Trapping by the positrons train: transfer matrix formulation*

$$\begin{pmatrix} \delta y_i(T=0) \\ \delta \dot{y}_i(T=0) \end{pmatrix} = \begin{bmatrix} \cosh(\Omega_p T) & \Omega_p^{-1} \sinh(\Omega_p T) \\ \Omega_p \sinh(\Omega_p T) & \cosh(\Omega_p T) \end{bmatrix} \begin{pmatrix} \delta y_i(0) \\ \delta \dot{y}_i(0) \end{pmatrix}$$

$$\begin{pmatrix} \delta y_i(t=T+0) \\ \delta \dot{y}_i(t=T+0) \end{pmatrix} = \begin{bmatrix} 1 & 0 \\ -\Omega_{K,\nu} & 1 \end{bmatrix} \begin{pmatrix} \delta y_i(t=T-0) \\ \delta \dot{y}_i(t=T-0) \end{pmatrix}$$

Eigen-value

$$\begin{bmatrix} \cosh(\Omega_p T) & \Omega_p^{-1} \sinh(\Omega_p T) \\ \Omega_p \sinh(\Omega_p T) & \cosh(\Omega_p T) \end{bmatrix} \begin{bmatrix} 1 & 0 \\ -\Omega_{K,\nu} & 1 \end{bmatrix} = e^{j\omega T} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

✓ *Dispersion relation*

Kick

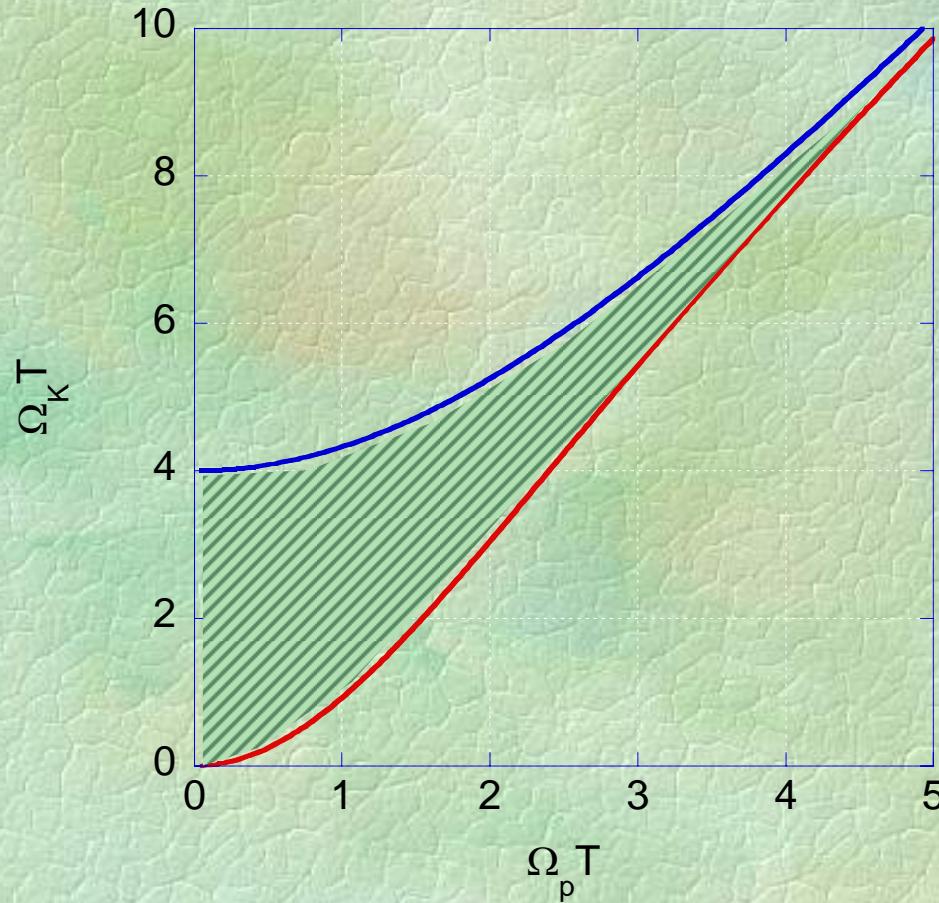
T- determines the scale

$$\cos(\omega T) = \cosh(\Omega_p T) - (\Omega_{K,\nu} T / 2\Omega_p T) \sinh(\Omega_p T)$$

Life-time

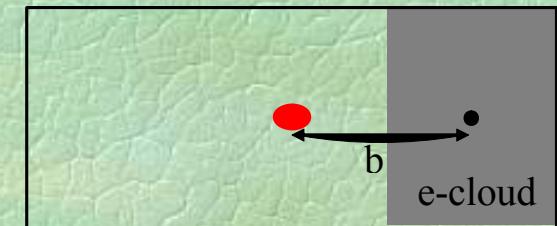
✓ *Trapping by the positrons train: Dispersion*

$$|\cos(\omega T)| < 1 \Rightarrow 2\Omega_p T \tanh(\Omega_p T / 2) < \Omega_{K,v} T < \frac{2\Omega_p T}{\tanh(\Omega_p T / 2)}$$



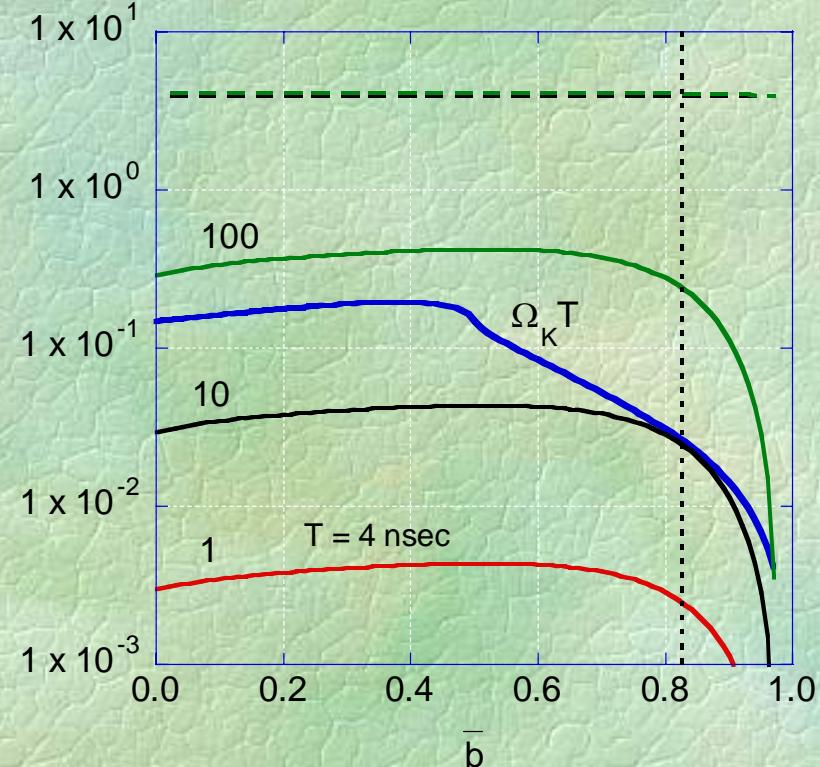
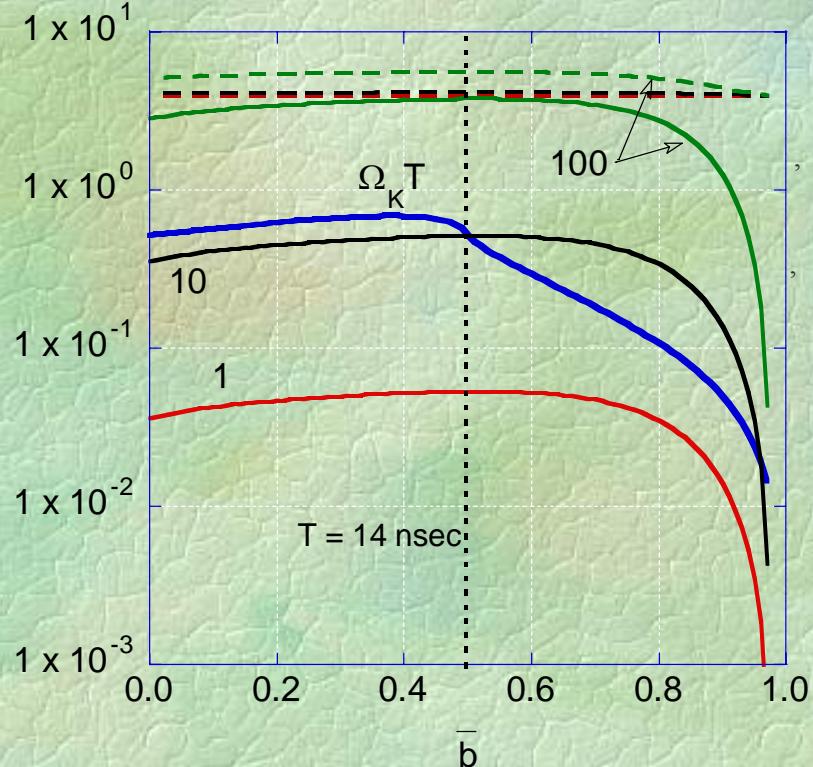
- ✓ For low values trapping is feasible
- ✓ At high values the constraint is very stringent

Life-time



✓ *Trapping by the positrons train: Dispersion*

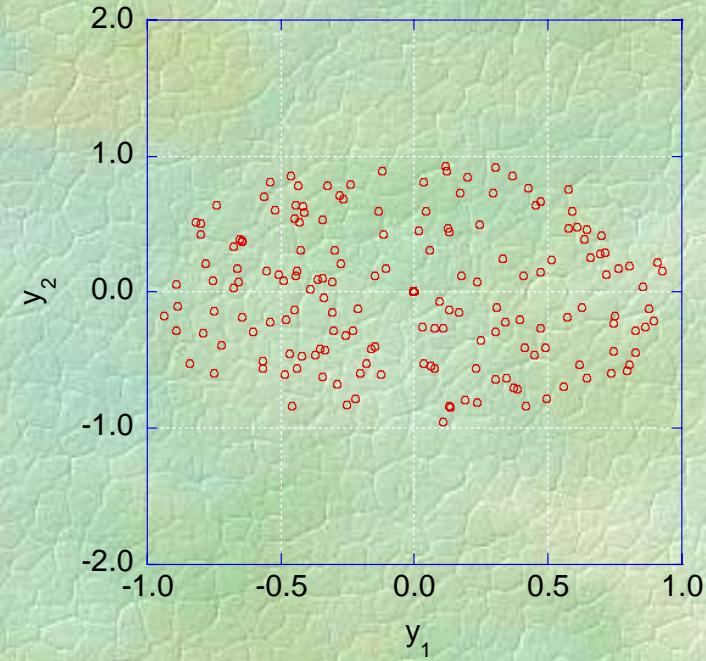
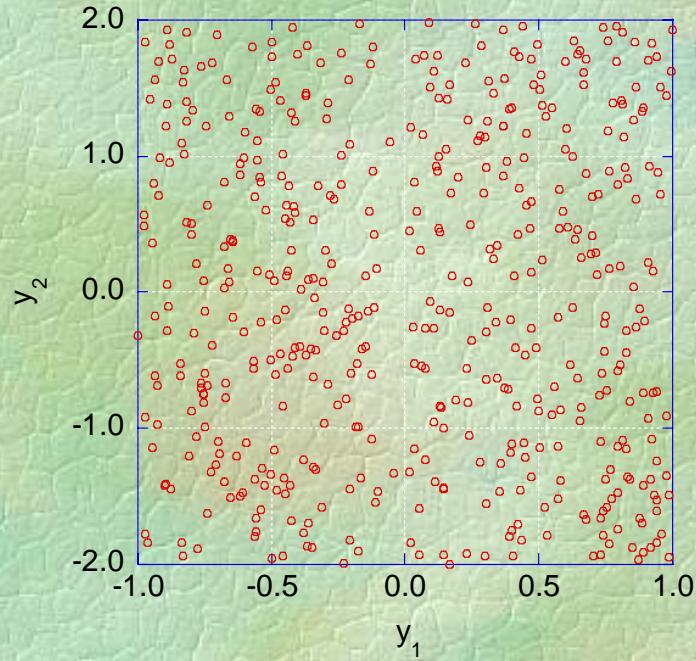
$$|\cos(\omega T)| < 1 \Rightarrow 2\Omega_p T \tanh(\Omega_p T / 2) < \Omega_{K,v} T < \frac{2\Omega_p T}{\tanh(\Omega_p T / 2)}$$



$$N_{pos} = 10^{10} \quad n_{ec} [m^{-3}] \times 10^{-11} = 1, 10 \text{ and } 100$$

Life-time

✓ *Trapping by the positrons train: Phase-space*



Only a fraction of the electrons are being trapped

$$\bar{y}_1^2 + \bar{y}_2^2 \leq 1$$

$$E \leq E_p [1 - \bar{y}_1^2] \leq E_p \equiv E_{\max,ps}$$

Life-time

Train of positron bunches

$$\tau = \underbrace{\int_0^{E_{\max,ps}} dE f_{ec}^{(p)}(E) \frac{1}{2} NT}_{\text{Trapped}} + \underbrace{\int_{E_{\max,ps}}^{\infty} dE f_{ec}^{(p)}(E) \tau'(E)}_{\text{Free}}$$

$$\simeq \underbrace{\frac{1}{2} NT \frac{E_{\max,ps}}{E_0}}_{\text{Trapped}} + \underbrace{\frac{1}{E_0 - E_{\max,ps}} \int_{E_{\max,ps}}^{E_0} dE \tau'(E) \left(1 - \frac{E_{\max,ps}}{E_0}\right)}_{\text{Free}}$$

Train of electron bunches

$$\tau = \int_0^{\infty} dE f_{ec}^{(e)}(E) \tau'(E) \simeq \frac{1}{E_0} \int_0^{E_0} dE \tau'(E)$$

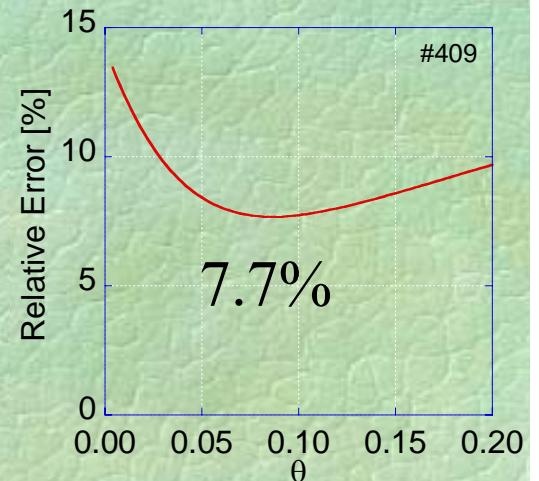
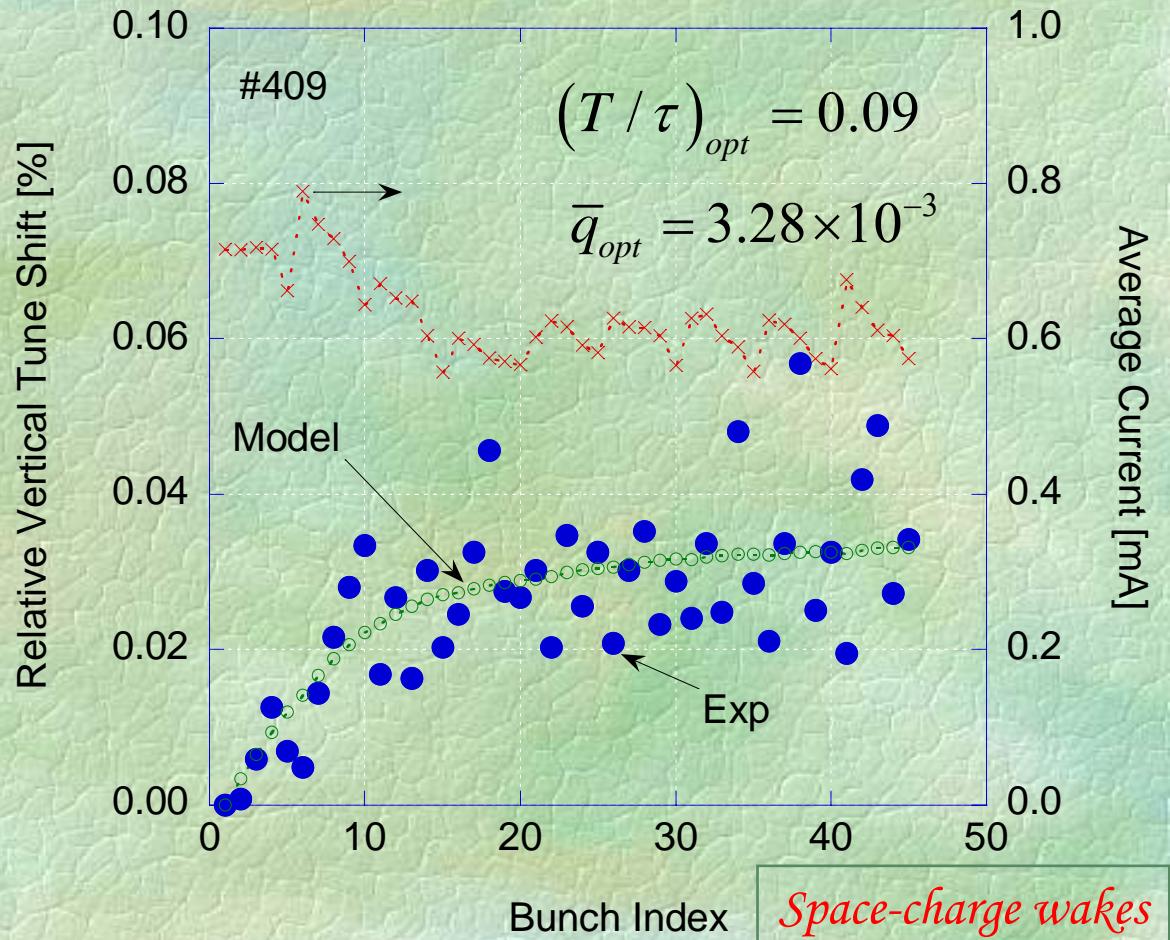
Theory and Experiment

$$\delta Q_\mu [\%] \approx \bar{q} \sum_{\nu=1}^N W_{\mu,\nu} \exp \left[-\frac{T}{\tau} (\mu - \nu) \right]$$

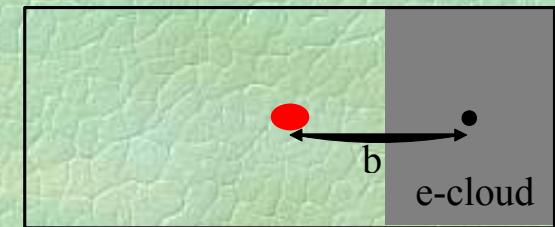
$$\bar{q} \equiv 50 \frac{\beta_r^2 r_e}{\gamma} \frac{D_x}{D_y} f_E(\bar{b}) n_{\text{nb}}$$

$$W_{\mu,\nu} = \frac{I_\nu}{\langle I \rangle} h(\mu - \nu - 0)$$

- ✓ Av. current $0.63[\text{mA}]$
- ✓ Vert. feedback=off
- ✓ Unknowns:
 $\langle E \rangle, b, \tau, n_{ec},$



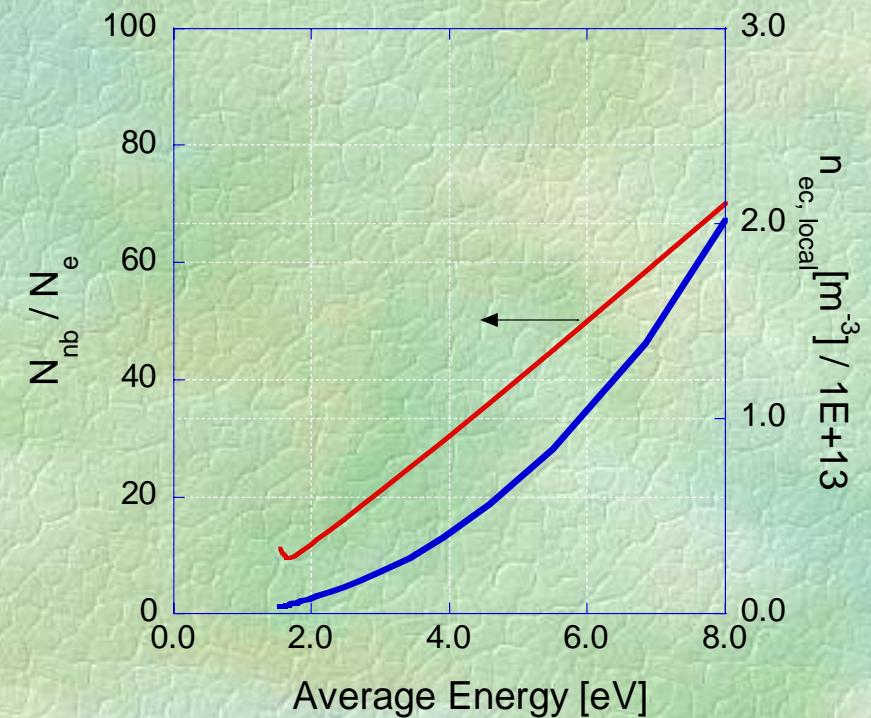
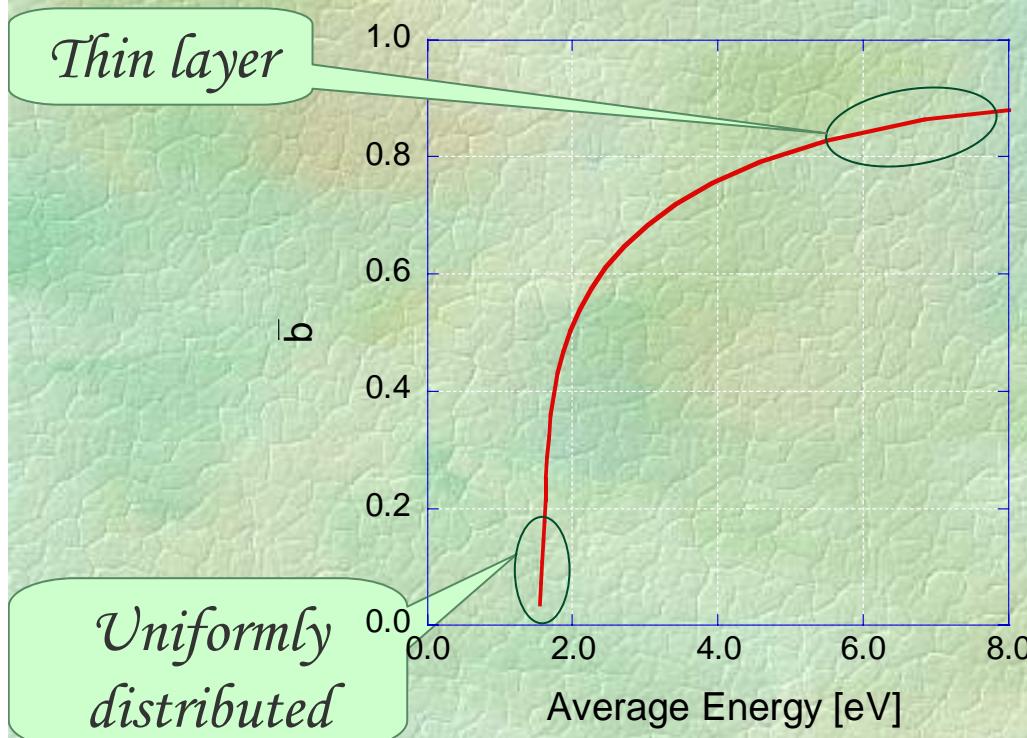
Theory and Experiment



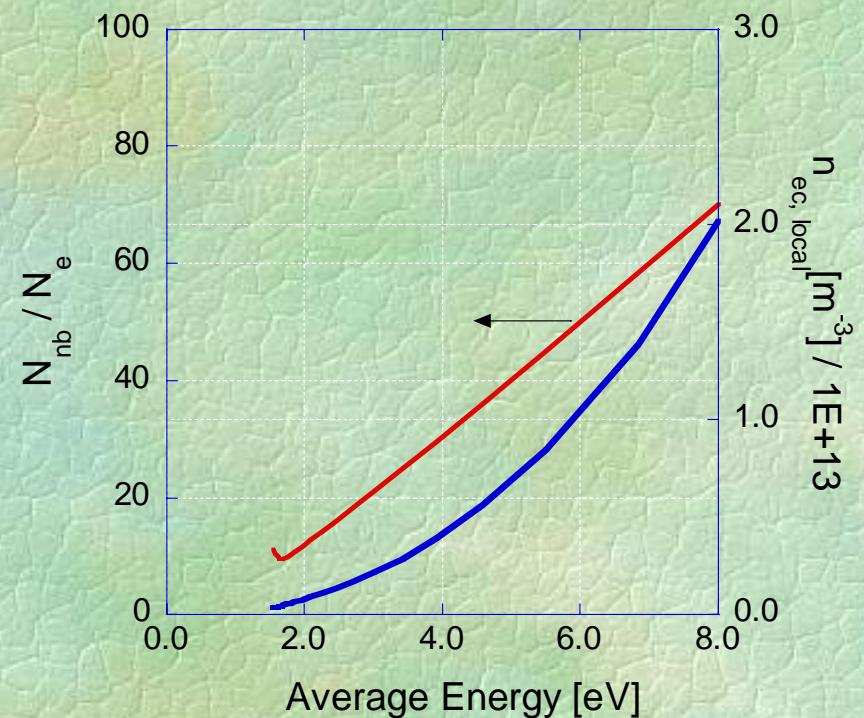
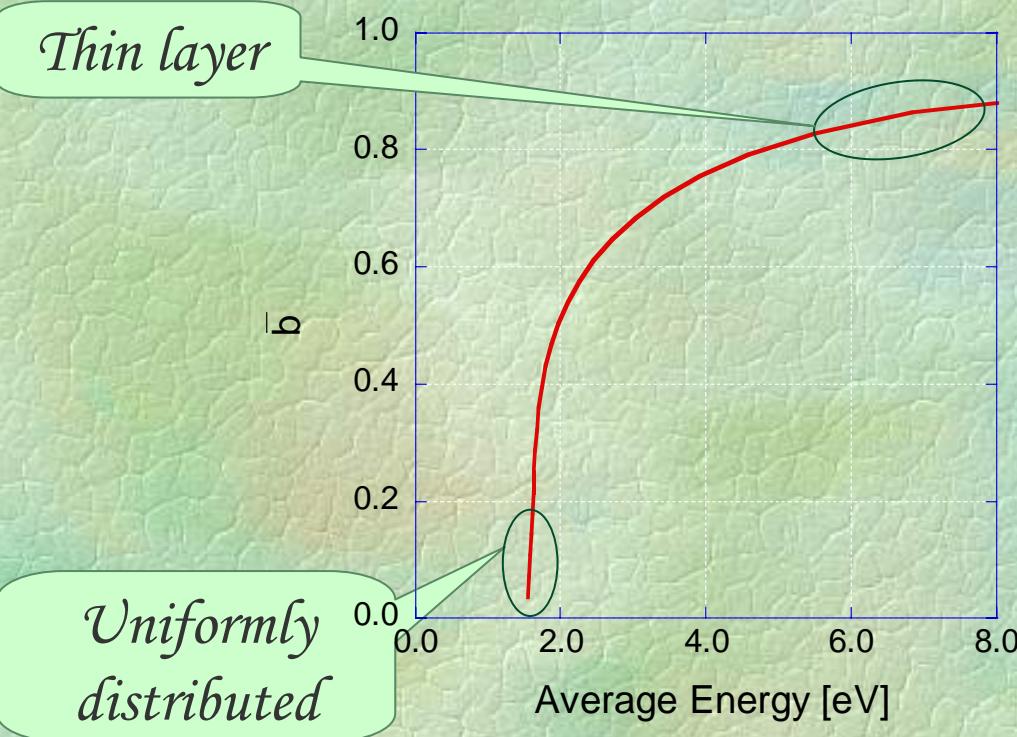
a) $(T / \tau)_{opt} = 0.09 \Rightarrow \tau_{opt} \simeq 155 \text{ n sec} \Rightarrow E_{eff} = 0.6 \text{ [eV]}$

b) $\bar{q}_{opt} = 3.28 \times 10^{-3} = 50 \frac{\beta_r^2 r_e}{\gamma} \frac{D_x}{D_y} f_E(\bar{b}) n_{nb} \Rightarrow f_E(\bar{b}) n_{nb} = 2 \times 10^{11} \text{ [m}^{-3}\text{]}$

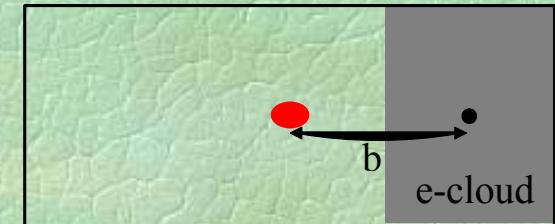
c) $n_{ec} \simeq n_{nb} \tau / T = n_{nb} / \theta_{opt} \simeq 11 n_{nb}$



Theory and Experiment



- ✓ With exception of the low-energy end, the model indicates that number of “new-borns” per positron, is **linear** with the average energy of the cloud
- ✓ Quadratic **local density** on the average energy ; less than 10^{11} e/m



Theory and Experiment

- ✓ The **photo-electrons** can account for only 1 new-born out of more than 10 (@2eV).

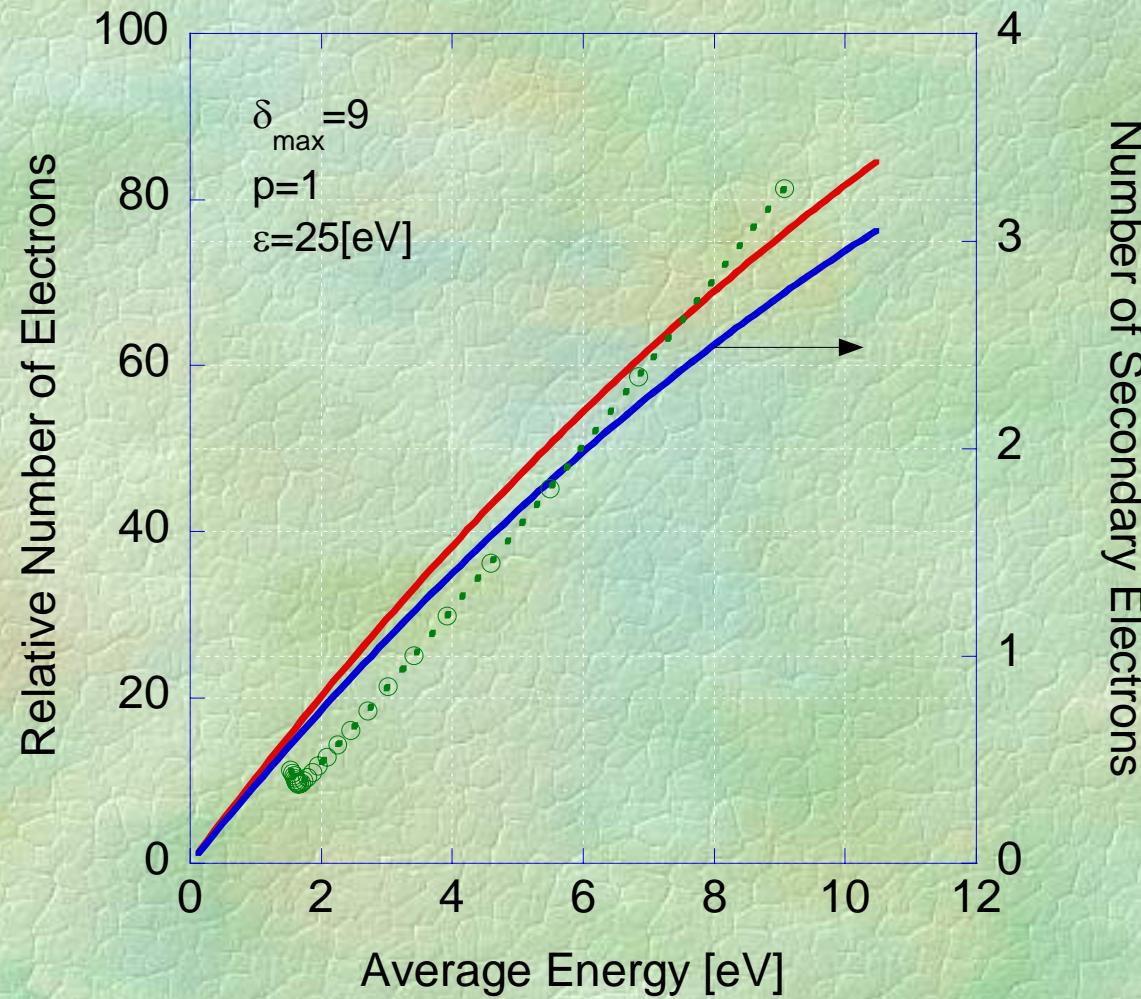
$$\frac{N_{\text{pe}}}{N_e} = 130 \times E_k [\text{GeV}] \bar{\delta}_{\text{pe}} \left(8 \frac{\langle E \rangle}{E_{\text{cr}}} \right) \sim 0.94 \text{@ 2eV}$$

- ✓ Secondary electrons contribute the remainder.

$$N_{\text{se}} = 130 \times E_k [\text{GeV}] \bar{\delta}_{\text{pe}} N_e \int_0^{\infty} dE f_{\text{se}}(E)$$

$$\langle E \rangle = \langle E \rangle_{\text{pe}} \frac{N_{\text{pe}}}{N_{\text{pe}} + N_{\text{se}}} + \langle E \rangle_{\text{se}} \frac{N_{\text{se}}}{N_{\text{pe}} + N_{\text{se}}}$$

Theory and Experiment



- ✓ If only SE are considered, the max SEY is high alluding to presence of Al_2O_3
- ✓ Or that other (linear) mechanisms play an important role ("stray" or bound)

Summary & Conclusions

- ✓ The **vertical tune-shift** is determined by the **average density** as experienced by each bunch
- ✓ **Build-up** of this average density is determined by two parameters:
“new-born” & life-time (“charging-up”: linear rise & equilibrium)
- ✓ “**New-born**” include:
 - photo electrons
 - secondary electrons
 - stray electrons
 - previously bound electrons (ionization)
- ✓ **Life-time** in case of a train of positrons has two contributions. Of:
 - fast electrons ($E > E_p$) that traverse the chamber
 - space-charge of the cloud
 - image-charge at the walls
 - trapped electrons ($E < E_p$) resembling ion trapping

Summary & Conclusions

- ✓ **Least-mean square fit** (model & experimental data) provides an estimate of the cloud's parameters: density, geometry and life-time.
- ✓ The **long life-time** of the cloud is not consistent with a multipactoring which requires energies of the order of 300[eV] or higher.
- ✓ Average energy of electrons in the cloud is a few electron volts.
- ✓ **Space-charge** waves may develop along this cloud and explain the “fluctuations” when the cloud reaches equilibrium

Acknowledgement

- ✓ *Introducing me to the subject: Maury Tigner, Mark Palmer & David Rice*
- ✓ *Great support with elucidation of the experimental data: Jerry Codner, Eugene Tanke and Robert Holtzapple*
- ✓ *Jim Crittenden was very helpful providing me with the detailed lattice data with and w/o wigglers.*
- ✓ *Kathey Harkay and John Flanagan for their input re this model.*
- ✓ *DoE and NSF*