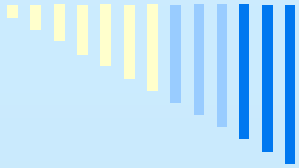


e-Cloud Theoretical Studies

Levi Schächter



Outline

- *Photo-Electrons*
- *Dynamics of Electrons in e-Cloud*
- *e-Cloud Spatial Distribution*
- *Wake in the e-Cloud*
- *Experimental Implications*

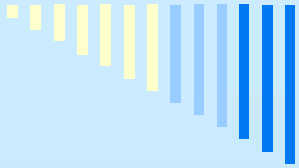


Photo-Electrons: synchrotron radiation

J.D. Jackson

$$\frac{N_{\text{ph}}(E > E_{\text{th}})}{N_e} \approx 49.54 \times E_k [\text{GeV}] \int_{E_{\text{th}}/E_{\text{cr}}}^{\infty} dy \int_{2y}^{\infty} dx K_{5/3}(x)$$

$$V_{\text{ph}}(E > E_{\text{th}}) = \frac{\mathcal{E}_{\text{ph}}(c/2\pi\rho)}{eN_e(c/2\pi\rho)} \approx 2.193 \times 10^5 \frac{E_k^4 (eV)}{\rho [m]} \int_{E_{\text{th}}/E_{\text{cr}}}^{\infty} dy y \int_{2y}^{\infty} dx K_{5/3}(x)$$

	$E_{\text{th}} [eV]$	$E_k = 2 [GeV]$	$E_k = 5 [GeV]$
N_{ph} / N_e	0	259.4	648.5
	4	177 (68%)	565 (87%)
	300	13.8 (5%)	314 (48%)
V_{ph}	0	177 [kV]	0.690 [MV]
	4	176 [kV]	0.690 [MV]
	300	6.6 [kV] (4%)	0.667 [MV] (97%)

Peak of SE (Al)

Work function (Al)

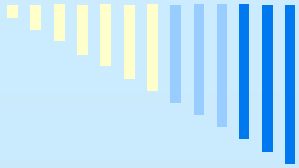


Photo-Electrons Yield

Photo-electrons *spectrum* is a “convolution” with the photon’s spectrum

$$\tilde{N}_{\text{pe}}(E) \approx \int_{E_w}^{\infty} dE' \delta_{\text{pe}}(E|E') [N_{\text{ph}} f_{\text{ph}}(E')]$$

Similar to the photons we are interested on :

$$\frac{\bar{N}_{\text{pe}}(E > E_{\text{th}})}{N_e} = 49.6 \times E_k [\text{GeV}] \int_{E_{\text{th}}}^{\infty} dE \int_{E_w}^{\infty} dE' \frac{1}{E_{\text{cr}}} \delta_{\text{pe}}(E|E') \int_{2E'/E_{\text{cr}}}^{\infty} dx K_{5/3}(x)$$

$$V_{\text{pe}}(E > E_{\text{th}}) = 2.193 \times 10^5 \frac{E_k^4 (\text{GeV})}{\rho [\text{m}]} \int_{E_{\text{th}}}^{\infty} dE \frac{E}{E_{\text{cr}}} \int_{E_w}^{\infty} dE' \delta_{\text{pe}}(E|E') \int_{2E'/E_{\text{cr}}}^{\infty} dx K_{5/3}(x)$$

Photo-Electrons – Model for the Yield

1. According to Grobner et.al [J. Vac. Sci. Technolo. A 7 (2) 1989 p, 223-9 “p.224 (Figure 2)] the yield of extracting a photo-electron from an *aluminium* surface

$$\delta_{\text{pe}}(E|E') \propto Y(E') = \left(\frac{E_0}{E'}\right)^\nu \quad \text{Typically: } \begin{cases} E_0 = 2.316[\text{eV}] \\ \nu = 0.825 \end{cases}$$

2. The energy of the photo-electron should not exceed the energy of the incident photon and we further assume that the electrons' energy is *uniformly distributed* between zero and the energy of incoming photon

$$\delta_{\text{pe}}(E|E') \propto \frac{1}{E'} [h(E) - h(E - E')]$$

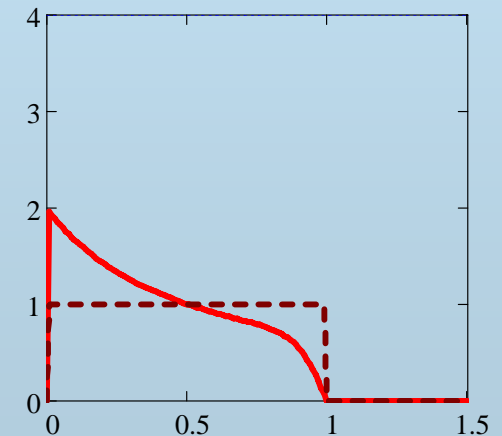




Photo-Electrons – Model for the Yield

3. No contribution from photons of energy lower than the *work function*

$$\delta_{\text{pe}}(E|E') \propto h(E' - E_w)$$

4. Not all the photons are absorbed

Reflection coefficient $\rho(E')$

$$\delta_{\text{pe}}(E|E') \propto \left[1 - |\rho_{\text{eff}}(E')|^2\right]$$

“Handbook of X-rays”, edited by E.F. Kaelble,
McGraw-Hill, New York, 1967, p. 48-2

$$|\rho_{\text{eff}}(E)|^2 \simeq \frac{1.881}{E} \exp(-1.2486 \times 10^{-4} E^2)$$

Virtually all photons with energy higher than 100eV are absorbed at the first impact!!

$$\delta_{\text{pe}}(E|E') \simeq Y(E') \frac{1}{E'} \left[h(E) - h(E - E') \right] h(E' - E_w) \left[1 - |\rho_{\text{eff}}(E')|^2 \right]$$

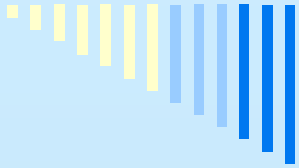


Photo-Electrons – Preliminary Conclusions

Comparable

	$E_{th} [eV]$	$E_k = 2 [GeV]$		$E_k = 5 [GeV]$	
		$\rho_{eff} (E) = 0$	$\rho_{eff} (E) \neq 0$	$\rho_{eff} (E) = 0$	$\rho_{eff} (E) \neq 0$
N_{pe} / N_e	0	25.8	20.9	32.9	27.6
	300	0.047(0.18%)	0.047(0.22%)	1.1(3.3%)	1.1(4%)
V_{pe}	0	354.4[V]	332.9[V]	1.64[kV]	1.61[kV]
	300	19.8[V](5.6%)	19.8[V](6%)	0.97[kV](59%)	0.97[kV](60%)
$\bar{\delta}_{pe}$	0	0.1(0)	0.8(-1)	0.5(-1)	0.4(-1)
	300	1.8(-4)	1.8(-4)	1.7(-4)	1.7(-4)

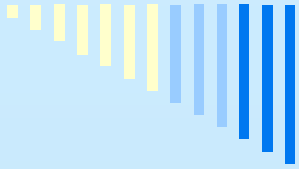
One order of magnitude difference

Reflections (Al) may reduce the number of photo-electrons by about 20%.



Outline

- *Photo-Electrons*
- *Dynamics of Electrons in e-Cloud*
- *e-Cloud Spatial Distribution*
- *Wake in the e-Cloud*
- *Experimental Implications*



Dynamics of Electrons in e-Cloud

$$\frac{\partial \tilde{N}_{ec}(t, E)}{\partial t} = \frac{1}{\tau(E)} \left[-\tilde{N}_{ec}(t, E) + \tilde{N}_{se}(t, E) \right] + \tilde{N}_{pe}(E) \sum_{n=1}^{N_b} \delta(t - T_n)$$

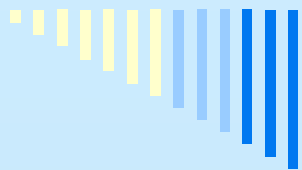
Life-time

Spectrum of SE

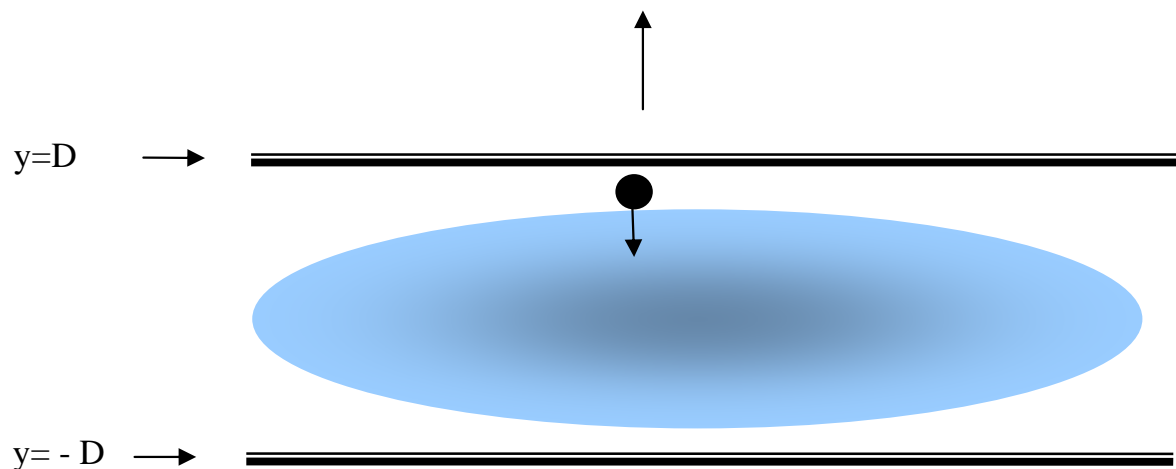
Spectrum of PE

*Bunch separation
is much longer than
the duration of
single bunch*

..or Ionization electrons

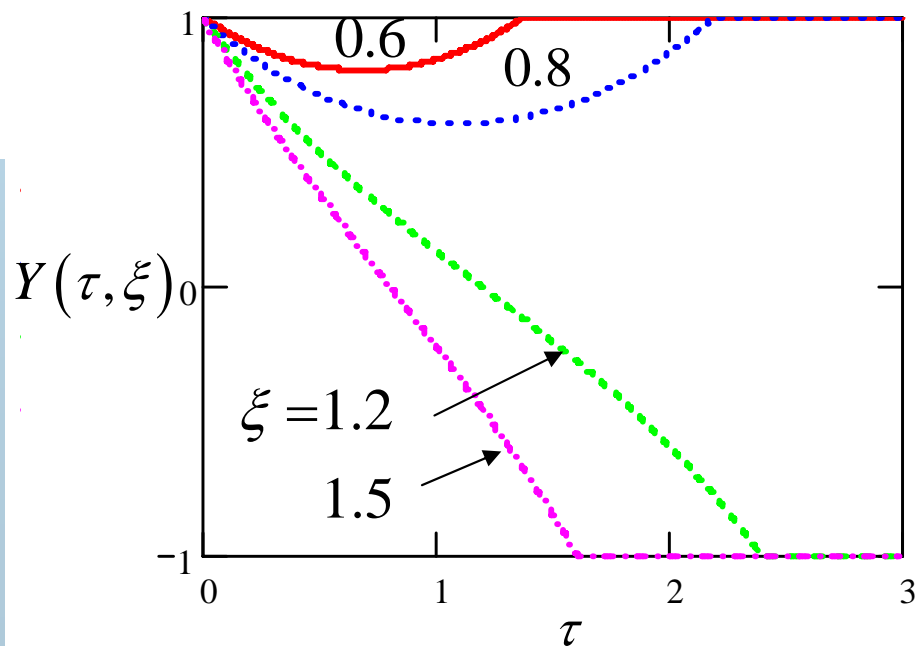


Dynamics of Electrons in e-Cloud $\tau(E) = ??$



$$\frac{d^2 y}{dt^2} = -\frac{e}{m} E_y = \frac{e^2 n_0}{m \epsilon_0} y = \omega_p^2 y$$

$$Y\left(\tau \equiv \omega_p t, \xi \equiv \frac{V_0}{\omega_p D}\right) \equiv \frac{y(t)}{D} = \cosh(\tau) - \xi \sinh(\tau)$$



Dynamics of Electrons in e-Cloud

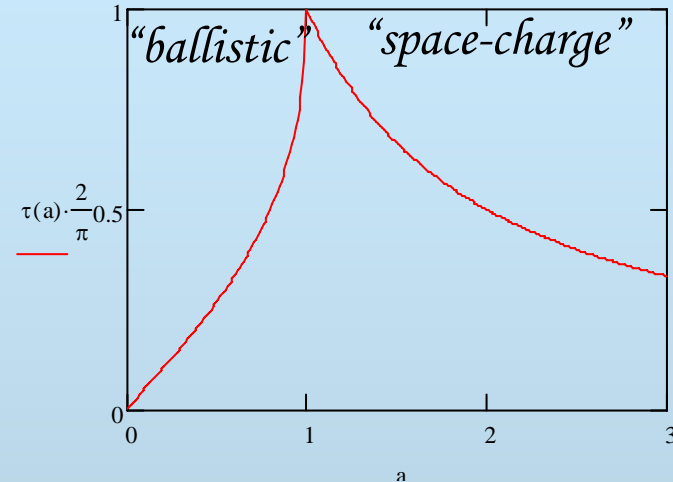
$$\tau(E) \approx \frac{\pi}{\omega_p} \begin{cases} \frac{\sqrt{E_0/E}}{2 - (E_0/E)^2} & E > E_0 \\ \sqrt{E/E_0} & E < E_0 \end{cases}$$

$$E_0 = \frac{1}{2} m \omega_p^2 a_y^2$$

$$E = E_0 = 100[\text{eV}], \quad D_y = 2.5[\text{cm}]$$

$$n_{\text{max}} [m^{-3}] = 1.1 \times 10^{12} \frac{E[\text{eV}]}{a_y^2[\text{cm}]} = 1.76 \times 10^{13}$$

$$\bar{\tau}^{(\text{max})} [n \text{ sec}] = 52.93 \frac{a_y[\text{cm}]}{\sqrt{E[\text{eV}]} } = 13.2$$



Note that the life-time of the electron in the cloud may be comparable to the bunch separation !!!

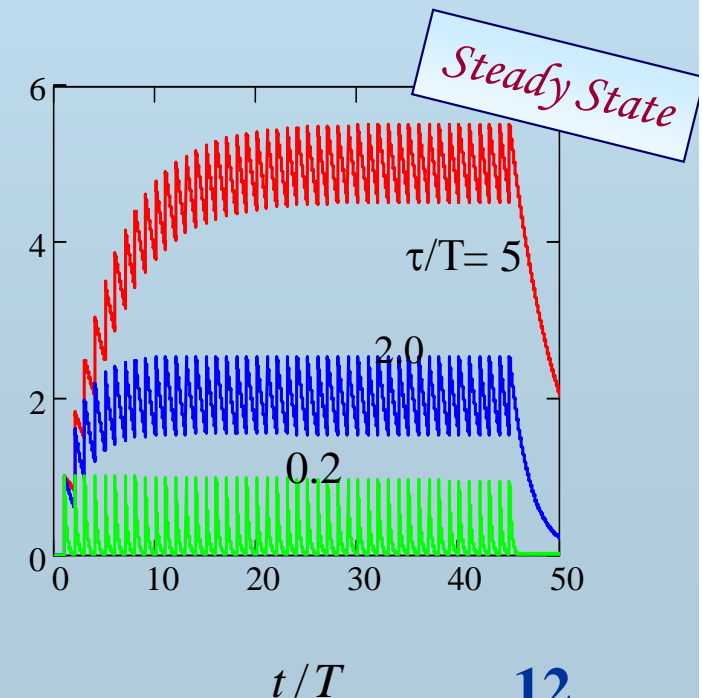
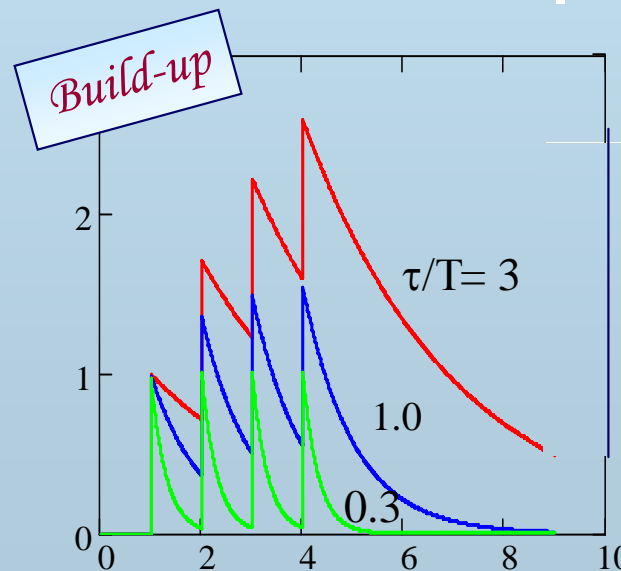
Dynamics of the e-Cloud

Let us consider the dynamics w/o SE

$$\frac{\partial \tilde{N}_{ec}(t, E)}{\partial t} = \frac{1}{\tau(E)} \left[-\tilde{N}_{ec}(t, E) + \tilde{N}_{se}(t, E) \right] + \tilde{N}_{pe}(E) \sum_{n=1}^{N_b} \delta(t - T_n)$$

$$\tilde{N}_{ec}(t, E) = \tilde{N}_{pe}(E) \sum_{n=1}^{N_b} h(t - T_n) \exp \left[-\frac{t - T_n}{\tau(E, \bar{n}_{ec})} \right]$$

Both build-up
and steady-state
depend on the
ratio T/τ



$$y_{n+1} = (y_n + \Delta) \exp(-T/\tau)$$

t/T

t/T

Dynamics of Electrons in e-Cloud

In practice, we need to integrate over the entire spectrum

$$\begin{aligned}\bar{n}_{ec}(t) &\equiv \frac{1}{(D_x D_y)(2\pi\rho)} \int_0^\infty dE \tilde{N}_{ec}(t, E) \\ &= \frac{\bar{N}_{pe}}{(D_x D_y)(2\pi\rho)} \sum_{n=1}^{N_b} h(t - T_n) \int_0^\infty dE f_{pe}(E) \exp\left[-\frac{t - T_n}{\tau(E, \bar{n}_{ec})}\right]\end{aligned}$$

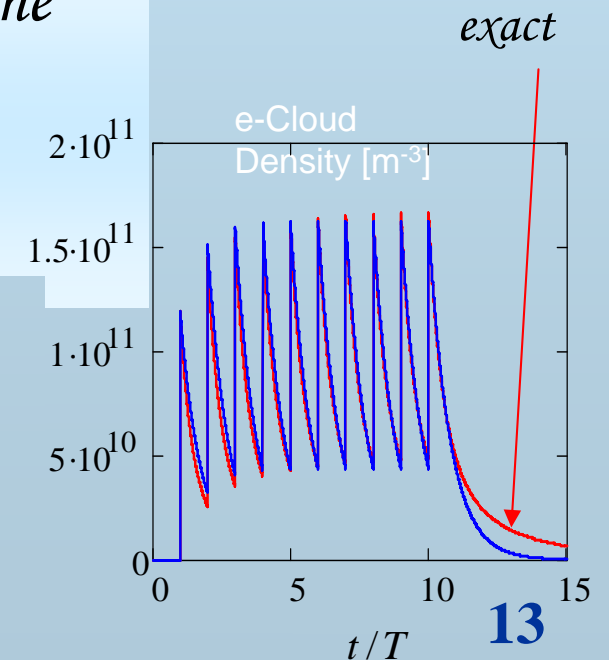
For the parameters involved, we found that defining the average life-time

$$\bar{\tau}(\bar{n}_{ec}) \equiv \int_0^\infty dE f_{pe}(E) \tau(E, \bar{n}_{ec})$$

is convenient since

$$\bar{n}_{ec}(t) = \frac{\bar{N}_{pe}}{(D_x D_y)(2\pi\rho)} \sum_{n=1}^{N_b} h(t - T_n) \exp\left[-\frac{t - T_n}{\bar{\tau}(\bar{n}_{ec})}\right]$$

No significant difference between the two



Dynamics of Electrons in e-Cloud

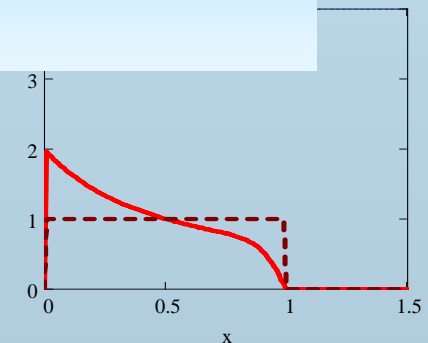
And now with *secondary emission* included:
$$\tilde{N}_{se}(t, E) = \int_0^{\infty} dE' \delta_{se}(E|E') \tilde{N}_{ec}(t, E')$$

Assumption: Because of the vertical magnetic field (bend or wiggler) the electrons gyrate along the magnetic field line consequently, both incoming and outgoing electrons are assumed to be moving in the *vertical* direction only.

Without significant loss of generality, based on the experimental data, it is reasonable to further assume that the energy of the secondary electrons is *uniformly distributed* between zero and that of the incident electron

$$\delta_{se}(E|E') \approx [h(E) - h(E - E')] \frac{\tilde{\delta}_{se}(E')}{E'}$$

Consequently,



$$\frac{\partial \tilde{N}_{ec}(t, E)}{\partial t} = \frac{-1}{\tau(E)} \left[\tilde{N}_{ec}(t, E) - \int_E^{\infty} dE' \frac{\tilde{\delta}_{se}(E')}{E'} \tilde{N}_{ec}(t, E') \right] + \tilde{N}_{pe}(E) \sum_{n=1}^{N_b} \delta(t - T_n)$$

Dynamics of Electrons in e-Cloud

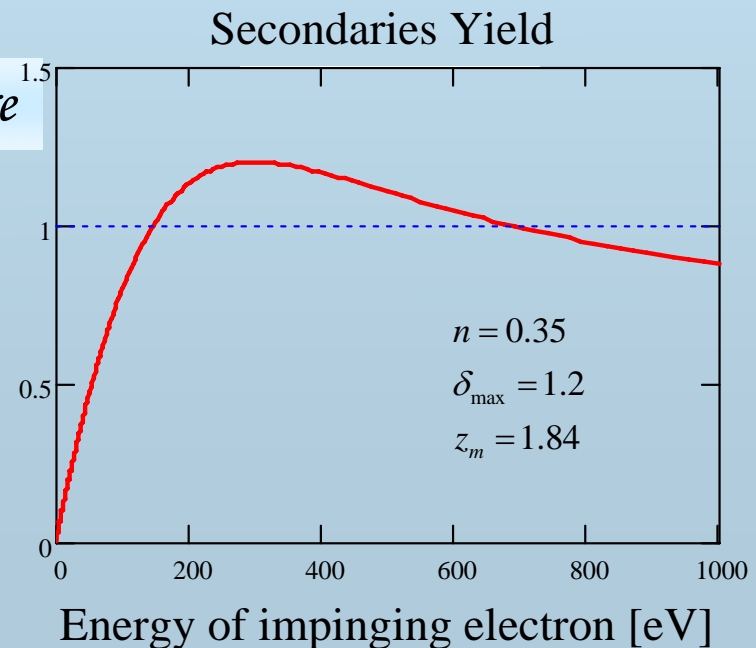
A realistic model for secondary emission [J.R.M Vaughan, "A new Formula for Secondary Emission Yield", IEEE Transactions on Electron Devices Vol 36 (9) p.1963 (1989). Actually a variant of: R.G. Lye and A.J. Dekker, "Theory of Secondary Emission", Physical Review Vol. 107, pp.977-981 (1957).]:

$$\tilde{\delta}_{se}(E') = \delta_{\max} \frac{1}{g_n(z_m)} g_n\left(z_m \frac{E}{E_{\max}}\right)$$

$$g_n(z) = \left[1 - \exp(-z^{n+1})\right] z^{-n}$$

Where the typical values of the two parameters are

$$n \approx 0.35, z_m \approx 1.84$$



Dynamics of Electrons in e-Cloud

Now we can proceed to a solution of

$$\frac{\partial \tilde{N}_{ec}(t, E)}{\partial t} = \frac{-1}{\tau(E)} \left[\tilde{N}_{ec}(t, E) - \int_E^{\infty} dE' \frac{\tilde{\delta}_{se}(E')}{E'} \tilde{N}_{ec}(t, E') \right] + \tilde{N}_{pe}(E) \sum_{n=1}^{N_b} \delta(t - T_n)$$

without SE we found

$$\tilde{N}_{ec}(t, E) = \tilde{N}_{pe}(E) \sum_{n=1}^{N_b} h(t - T_n) \exp\left[-\frac{t - T_n}{\tau(E)}\right]$$

with SE the solution is

$$\tilde{N}_{ec}(t, E) = \tilde{N}_{pe}(E) \sum_{n=1}^{N_b} h(t - T_n) \exp\left[-\frac{t - T_n}{\tau_{ec}(E)}\right]$$

wherein

$$\tau_{ec}(E) = \tau(E) + \frac{\exp[-\psi(E)]}{\tilde{N}_{pe}(E)} \int_E^{\infty} dE' \exp[\psi(E)] \frac{\tilde{\delta}_{se}(E')}{E'} \tilde{N}_{pe}(E') \tau(E')$$

$$\psi(E) = \int_0^E dE' \frac{\tilde{\delta}_{se}(E')}{E'}$$

Dynamics of Electrons in e-Cloud

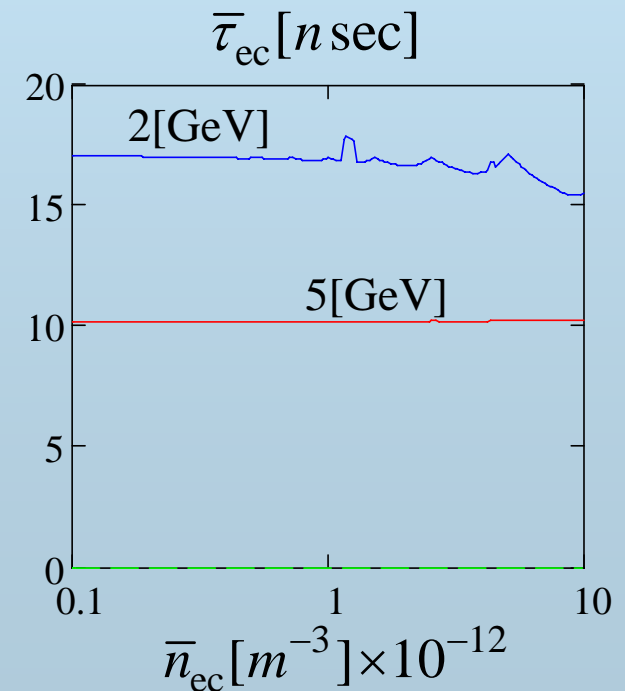
The average density during the train duration is

$$\begin{aligned}\bar{n}_{ec} &\equiv \frac{1}{(D_x D_y)(2\pi\rho)} \frac{1}{TN_b} \int_0^\infty dt \int_0^\infty dE \tilde{N}_{ec}(t, E) \\ &= \frac{\bar{N}_{pe}}{(D_x D_y)(2\pi\rho)} \frac{1}{T} \int_0^\infty dE f_{pe}(E) \tau_{ec}(E)\end{aligned}$$

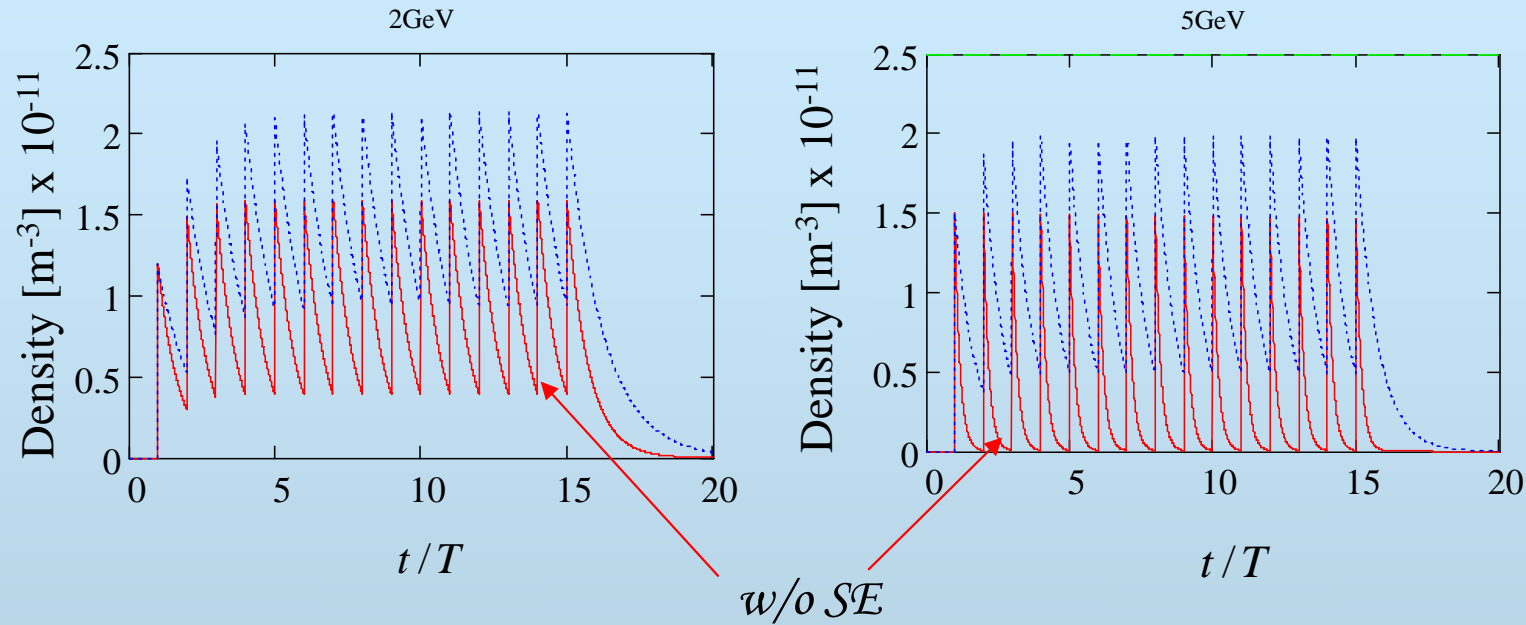
$$\bar{\tau}_{ec} \equiv \int_0^\infty dE f_{pe}(E) \tau_{ec}(E)$$

- For a 5GeV train, the average decay is 10nsec in comparison to about 2.5nsec when no SE were present.
- At 2GeV the average decay increases from 10 to 17nsec.
- In both cases the average decay is virtually independent of the density.

This result can be attributed to SE being significantly slower than the PE and as a result, it takes them longer time to return to the wall or traverse the gap



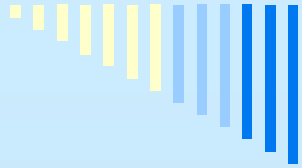
Dynamics of Electrons in e-Cloud



Relying on discharge analogy

$$N_{\text{ec}}^{(\text{low})} = \frac{\bar{N}_{\text{pe}}}{\exp(T / \bar{\tau}_{\text{ec}}) - 1}$$

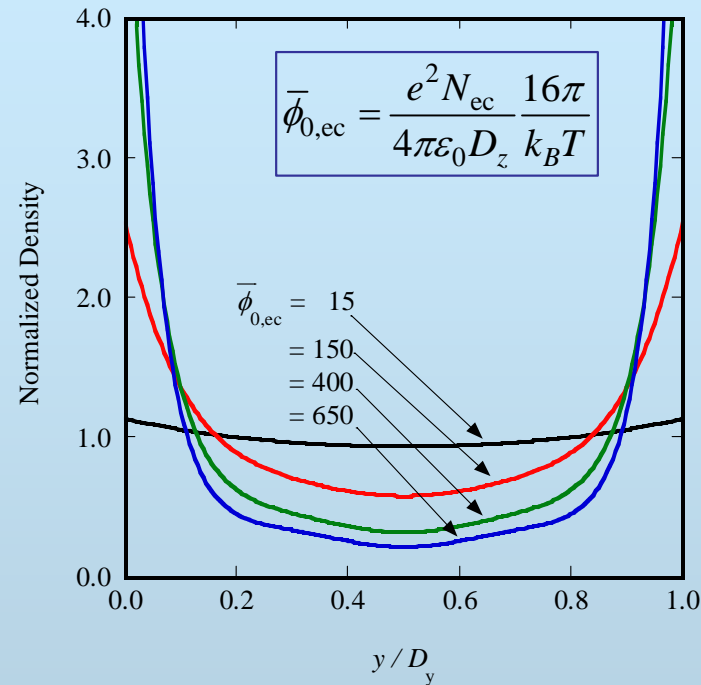
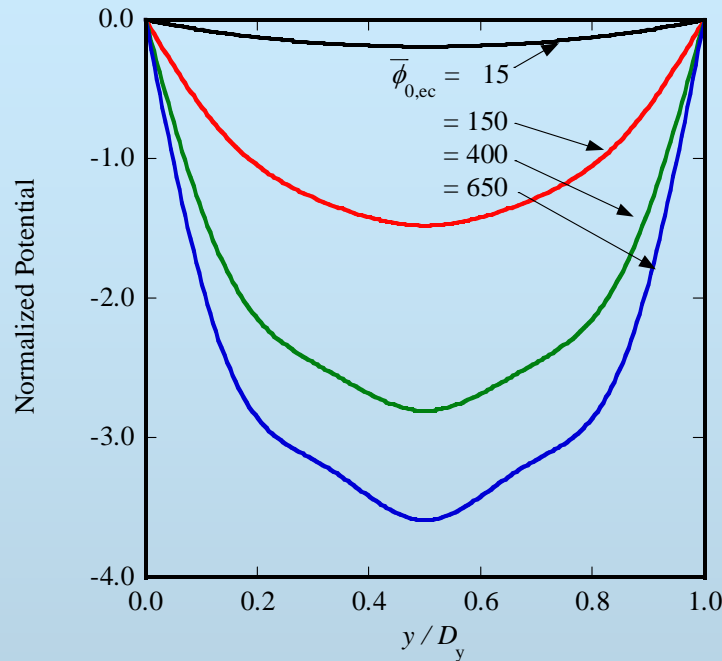
$$N_{\text{ec}}^{(\text{high})} \simeq \bar{N}_{\text{pe}} + \frac{\bar{N}_{\text{pe}}}{\exp(T / \bar{\tau}_{\text{ec}}) - 1} = \frac{\bar{N}_{\text{pe}}}{1 - \exp(-T / \bar{\tau}_{\text{ec}})}$$



Outline

- *Photo-Electrons*
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- *e-Cloud Static Spatial Distribution*
- *Wake in the e-Cloud*
- *Experimental Implications*

e-Cloud Static Spatial Distribution



Note that in any case, in the center, the potential well may be approximated by the potential of a *harmonic oscillator*.

$$\bar{\phi}(\bar{y}) \approx \bar{\phi}(\bar{y} = 1/2) + \cancel{\frac{d\bar{\phi}}{d\bar{y}} \Big|_{\bar{y}=1/2}} \delta\bar{y} + \frac{1}{2} \frac{d^2\bar{\phi}}{d\bar{y}^2} \Big|_{\bar{y}=1/2} \delta\bar{y}^2 \quad \kappa = \frac{d^2\bar{\phi}}{d\bar{y}^2} \Big|_{\bar{y}=1/2} = -\sum_{n=1}^{\infty} (\pi n)^2 \bar{\phi}_n \sin\left(\frac{\pi}{2}n\right)$$

Spring coefficient



e-Cloud Static Spatial Distribution

Implication on the vertical tune. Let the oscillatory motion (equilibrium) in the absence of the e-cloud be described by

$$\left(\frac{d^2}{dt^2} + \frac{c^2}{\beta_0^2} \right) \delta y = 0$$

In the presence of the e-cloud

$$\left(\frac{d^2}{dt^2} + \frac{c^2}{\beta_0^2} \right) \delta \bar{y} = \frac{-e}{m\gamma D_y} \left(-\frac{\partial \phi}{\partial y} \right) \approx \frac{5}{m\gamma D_y^2} \frac{e^2 N_{ec}}{4\pi\epsilon_0 D_z} \delta \bar{y}$$

Implying

$$\frac{1}{\beta^2} = \frac{1}{\beta_0^2} - \frac{5N_{ec}r_e}{\gamma D_y^2 D_z} = \frac{1}{\beta_0^2} - 5n_{ec} \frac{r_e}{\gamma} \frac{D_x}{D_y}$$

$$\frac{\Delta v}{v_0} = \sqrt{1 - 5 \frac{n_{ec} r_e \beta_0^2}{\gamma} \frac{D_x}{D_y}} - 1 \approx -\frac{5}{2} \frac{n_{ec} r_e \beta_0^2}{\gamma} \frac{D_x}{D_y} = -V n_{ec}$$

*During the passage of the bunch, temporarily, the distribution changes. Electrons **repel** the cloud and positron **attract** it making the potential shallower.*

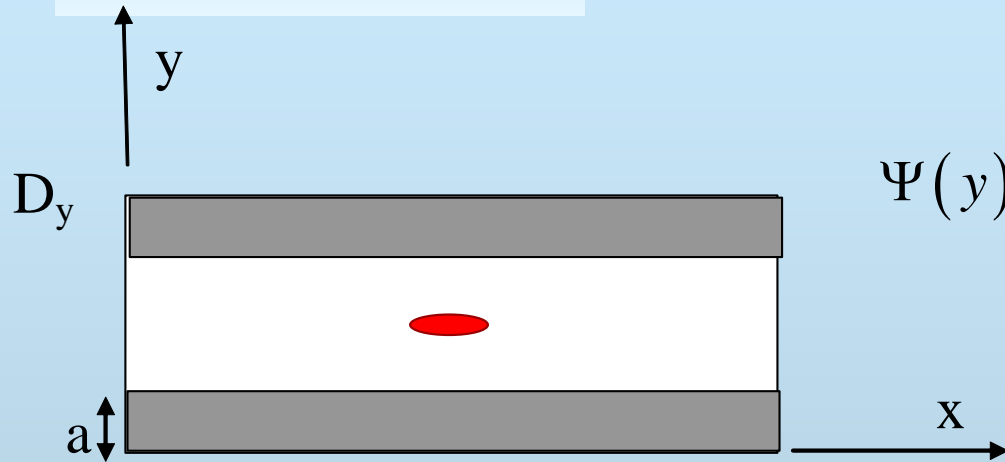


Outline

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Wakes in the e-Cloud

Eigen-frequencies



$$E_y(x, y, z) = \Psi(y) \sin\left(\pi n_x \frac{x}{D_x}\right) \exp\left(-j \frac{\omega}{V} z\right)$$

$$\Psi(y) = \begin{cases} A \cosh(\Gamma y) & 0 < y < a \\ B \sinh\left[\frac{\pi n_x}{D_x} \left(y - \frac{D_y}{2}\right)\right] & a < y < D_y/2 \end{cases}$$

$$\Gamma^2 = \left(\frac{\pi^2 n_x^2}{D_x^2} + \frac{\omega_p^2}{c^2}\right) \left(1 - \frac{\omega_p^2}{\omega^2}\right)^{-1}, \quad \Gamma a = j\xi$$

$$K_{n_x}^2 \equiv \left(\frac{\pi n_x a}{D_x}\right)^2 + \left(\frac{\omega_p}{c} a\right)^2$$

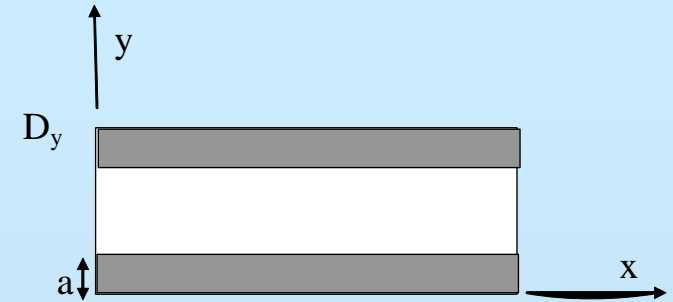
$$\xi_0 \equiv \left(\frac{\pi n_x a}{D_x}\right) \coth\left(\frac{\pi n_x a}{D_x} \left(\frac{D_y}{2a} - 1\right)\right)$$

$$\tan\left(\frac{\xi}{\xi_0}\right) = \frac{\xi_0}{\xi}$$

$$\omega_i^2 = \omega_p^2 \frac{\xi_i^2}{\xi_i^2 + K_{n_x}^2}$$

Wakes in the e-Cloud

Generated Power



$$P = \int dx \int dy \int dz \mathbf{J}_z \mathbf{E}_z$$

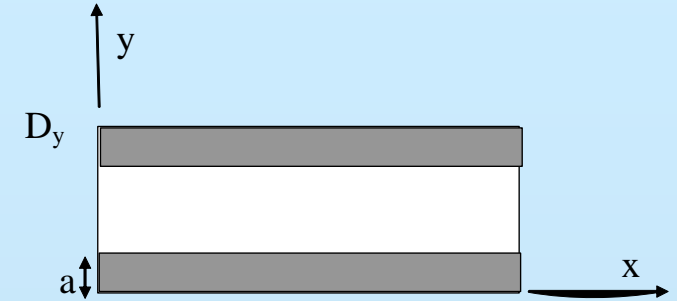
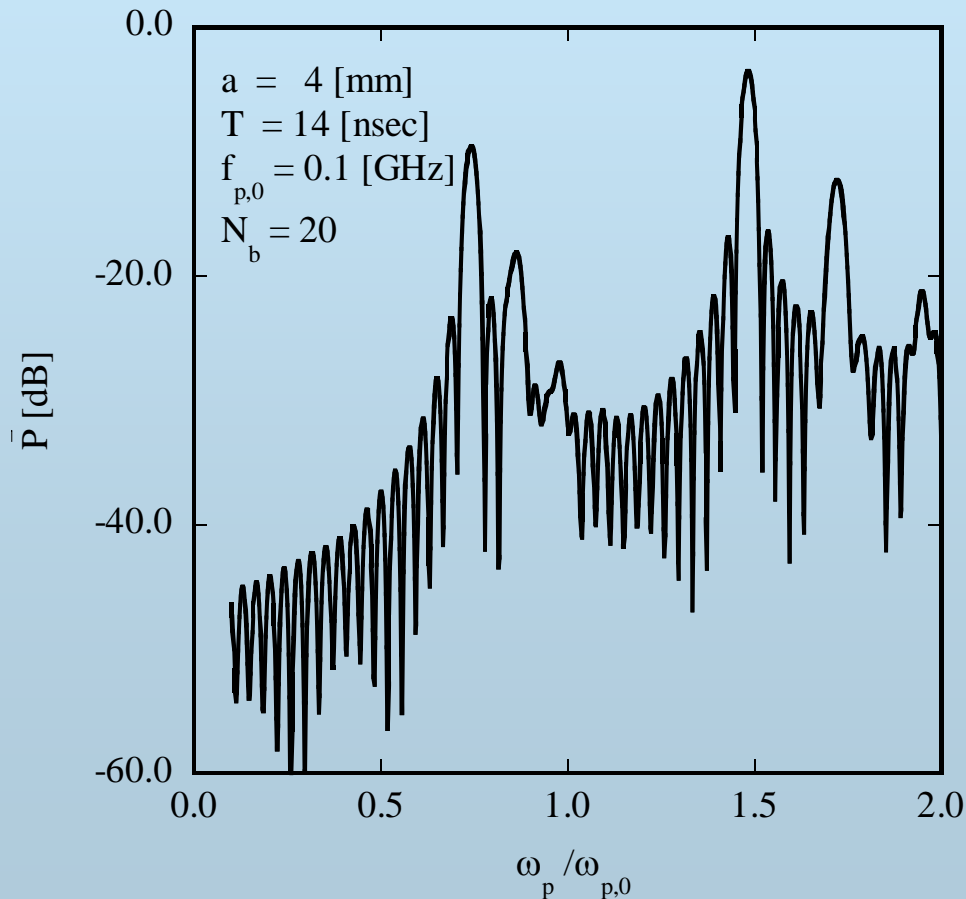
$$\bar{P} = \frac{-P}{\eta_0 \left(\frac{q_e c}{D_x} \right)^2} = 2 \left(\frac{a}{D_x} \right) \left(\frac{\omega_p}{c} D_x \right)^2 \sum_{n_x=0}^{\infty} \frac{\sin^2 \left(\frac{\pi}{2} n_x \right)}{\sinh^2 \left(\psi_{n_x} \right)} \left(\pi n_x \frac{a}{D_x} \right)^2$$

$$\times \sum_i \frac{\xi_i^2 \left[\xi_i^2 + \left(\pi n_x \frac{a}{D_x} \right)^2 \right]^{-1}}{\left[\xi_i^2 + \xi_0 (1 + \xi_0) \right] \left(\xi_i^2 + \kappa_{n_x}^2 \right)} \left[\frac{N_b^2 \operatorname{sinc}^2 \left(\frac{1}{2} \omega_i T N_b \right)}{2 \operatorname{sinc}^2 \left(\frac{1}{2} \omega_i T \right)} \right]$$

Eigen-frequencies

Wakes in the e-Cloud

Generated Power

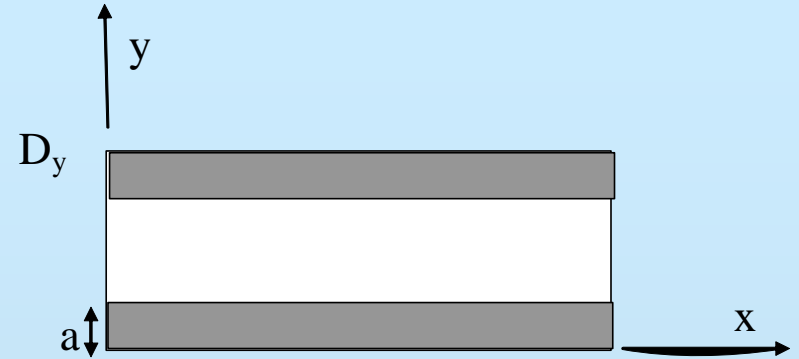


- *Cloud-density controls emitted power*
- *Longer train, sharper peaks*
- *At peak value (infinite life-time)*

$$P \sim -\eta_0 \left(eN_e c / D_x \right)^2 N_b^2 / 2$$

Wakes in the e-Cloud

Transverse Kick: “spring coefficient”



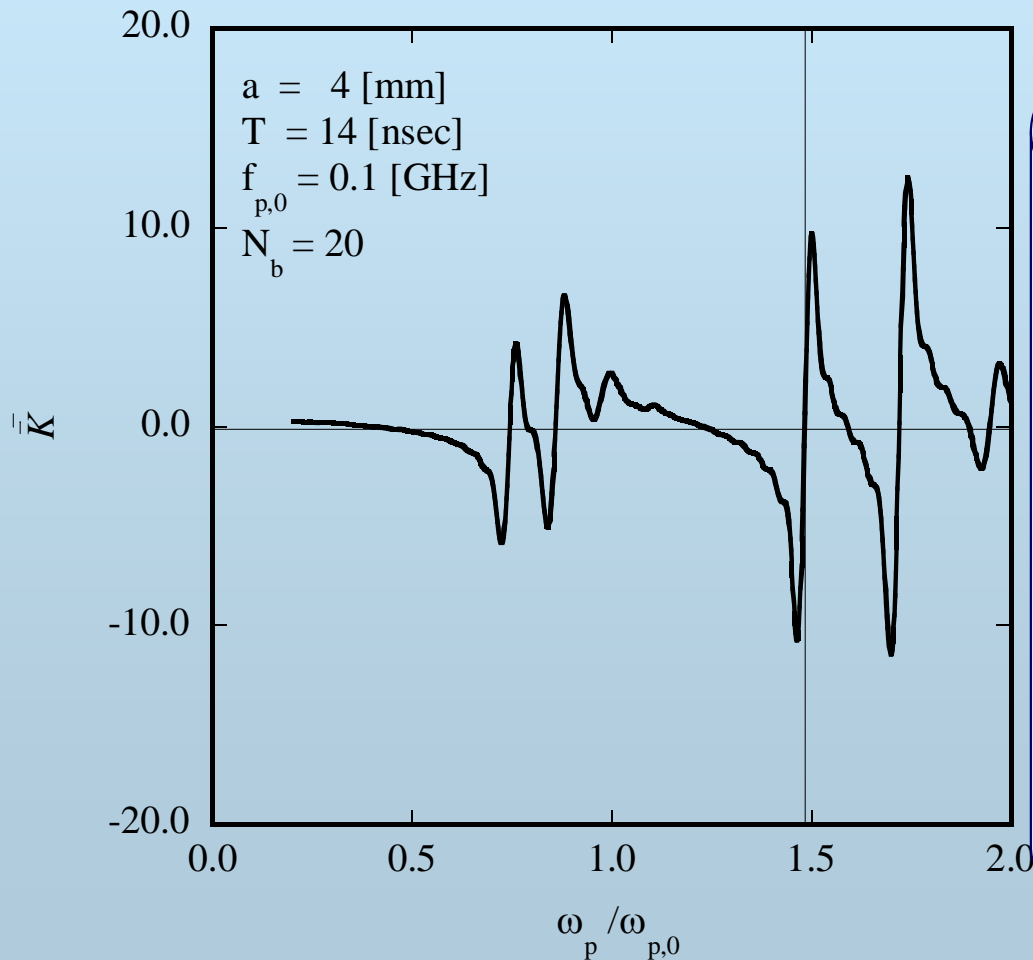
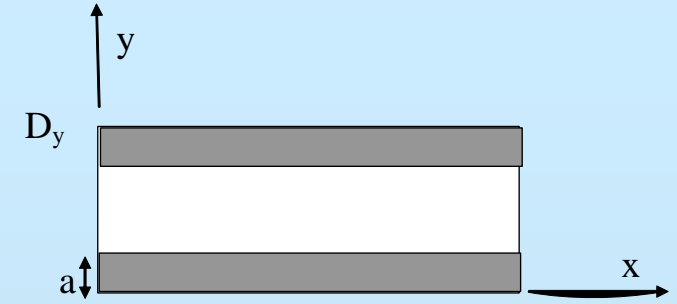
$$\bar{K} = \frac{f_y(y) D_x c}{\eta_0 \left(\frac{q_e c}{D_x} \right)^2 \left[g(y, \Delta_y) \left(y - \frac{D_y}{2} \right) \right]}$$

$$= 2 \left(\frac{\omega_p}{c} D_x \right)^2 \frac{Tc}{a} \sum_{n_x=0}^{\infty} \left(\pi n_x \frac{a}{D_x} \right)^4 \frac{\sin^2 \left(\pi n_x \frac{1}{2} \right)}{\sinh^2 \psi_{n_x}}$$

$$\times \sum_i \frac{\xi_i^2 \left[\xi_i^2 + \left(\frac{\pi n_x a}{D_x} \right)^2 \right]^{-1}}{\left[\xi_i^2 + \chi_{n_x} (1 + \chi_{n_x}) \right] (\xi_i^2 + \kappa_{n_x}^2)} N_b \frac{\text{sinc}(\omega_i T) - \text{sinc}(\omega_i T N_b)}{(\omega_i T)^2 \text{sinc}^2 \left(\frac{1}{2} \omega_i T \right)}$$

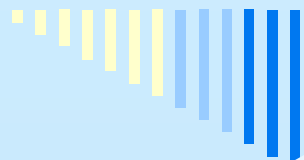
Wakes in the e-Cloud

Transverse Kick: “spring coefficient”



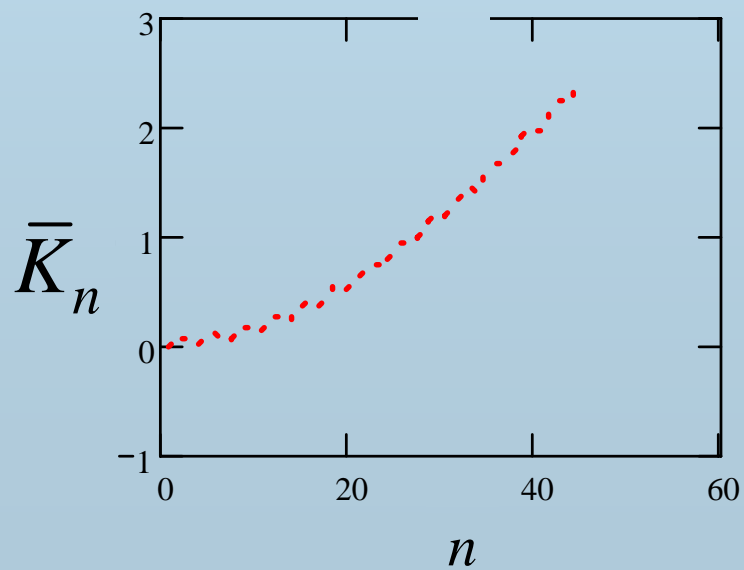
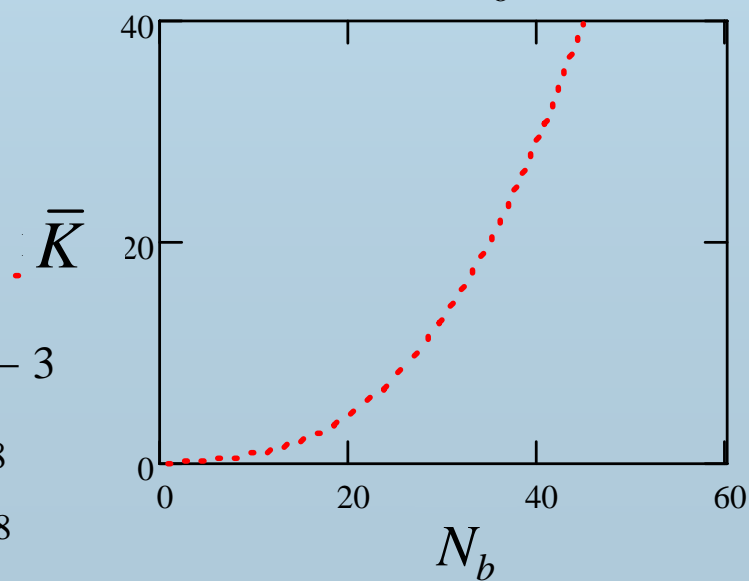
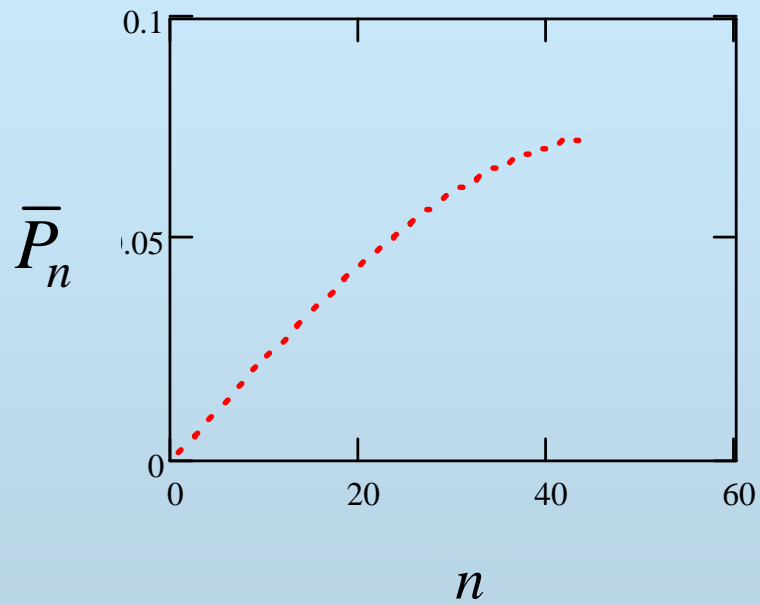
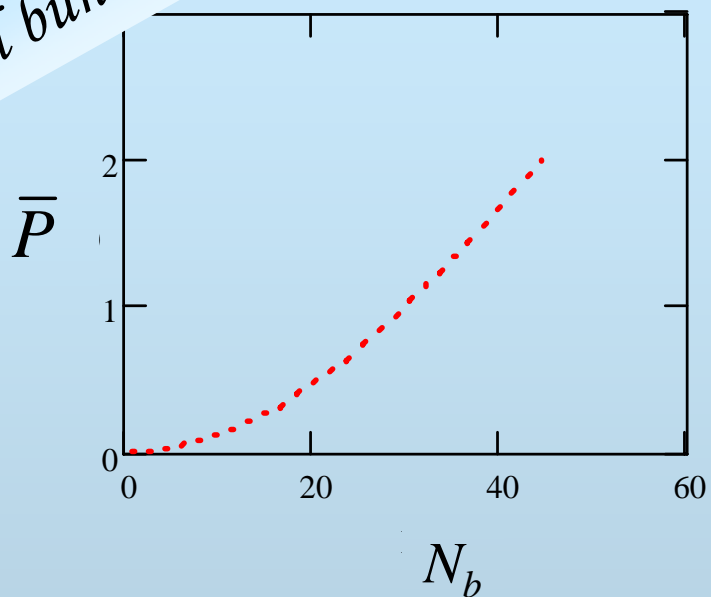
The spring coefficient as a function of the cloud density, may be positive in which case the vertical dynamics is unstable or it can become stable for densities that $K < 0$.

This behavior is affected by both the cloud density and the cloud thickness



Wakes in the e-Cloud

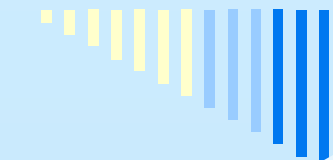
Effect on the individual bunch



$$a = 4.3 \times 10^{-3}$$

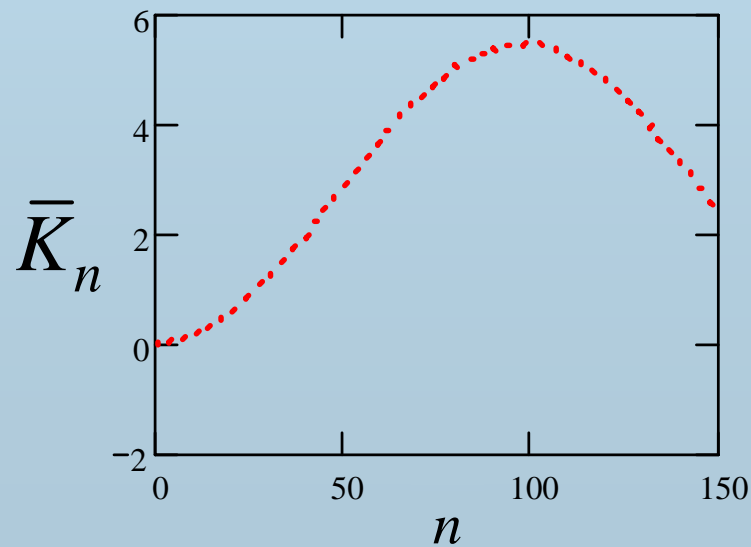
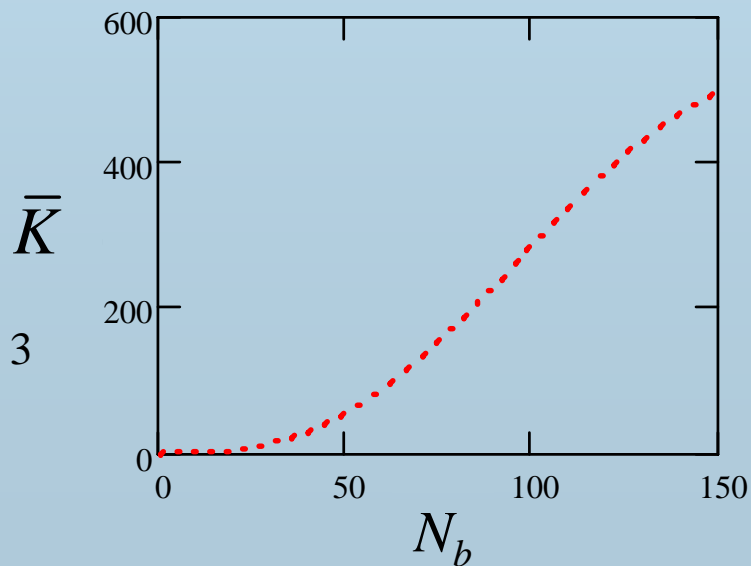
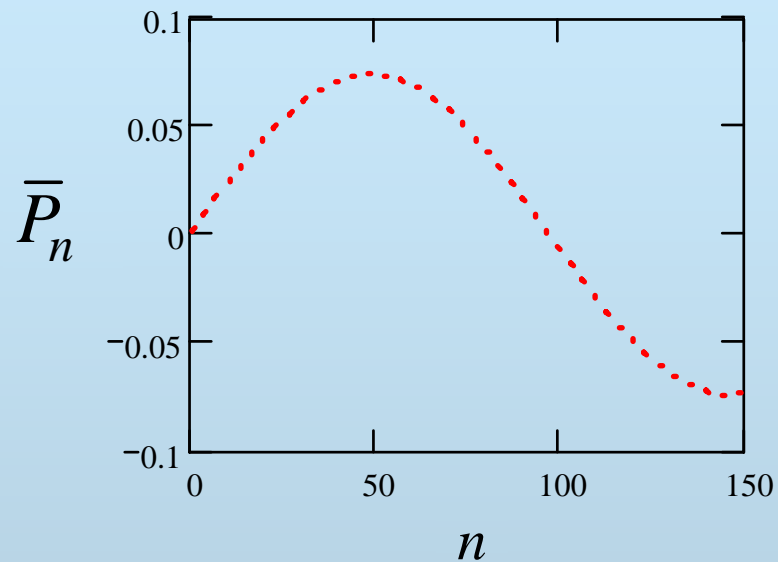
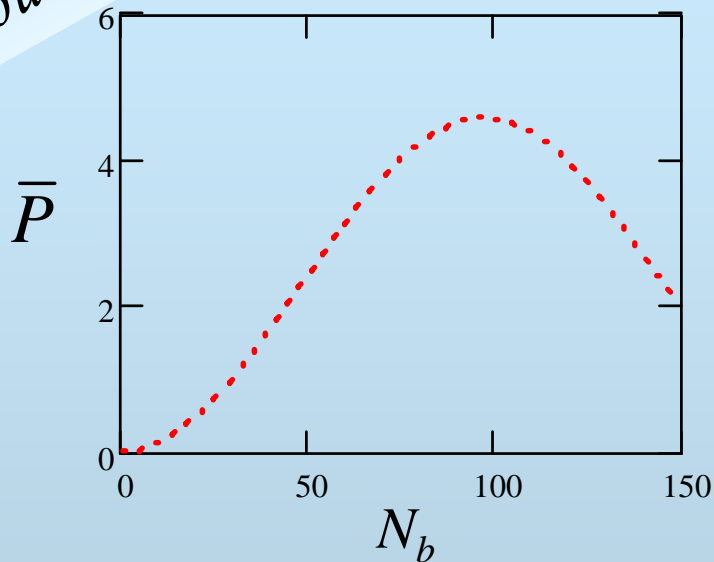
$$T = 1.4 \times 10^{-8}$$

$$f_p = 1.483 \times 10^8$$



Wakes in the e-Cloud

Effect on the individual bunch

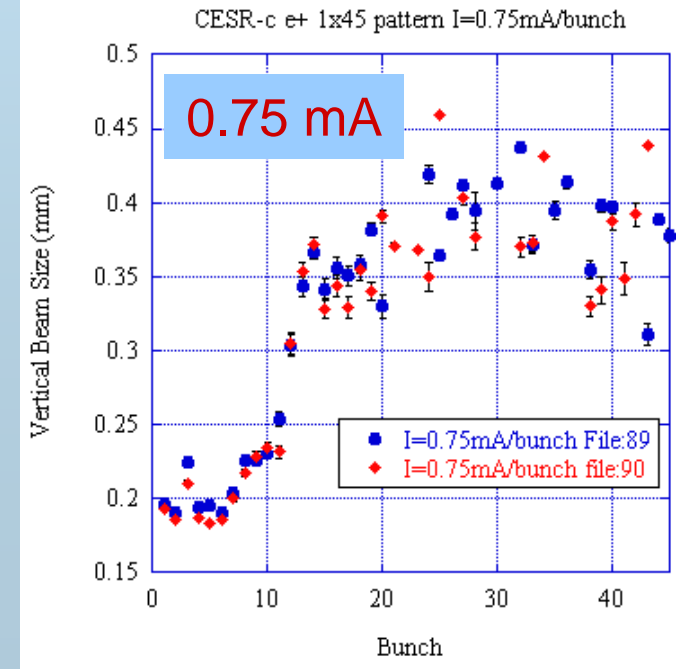
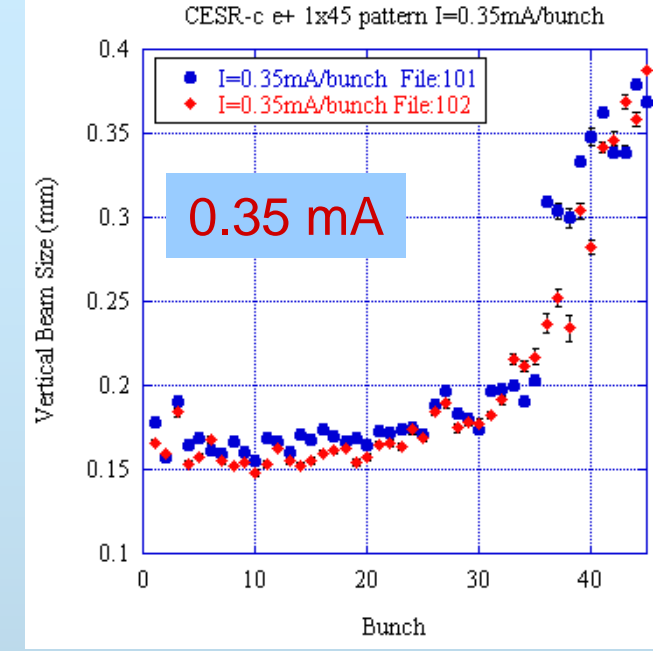
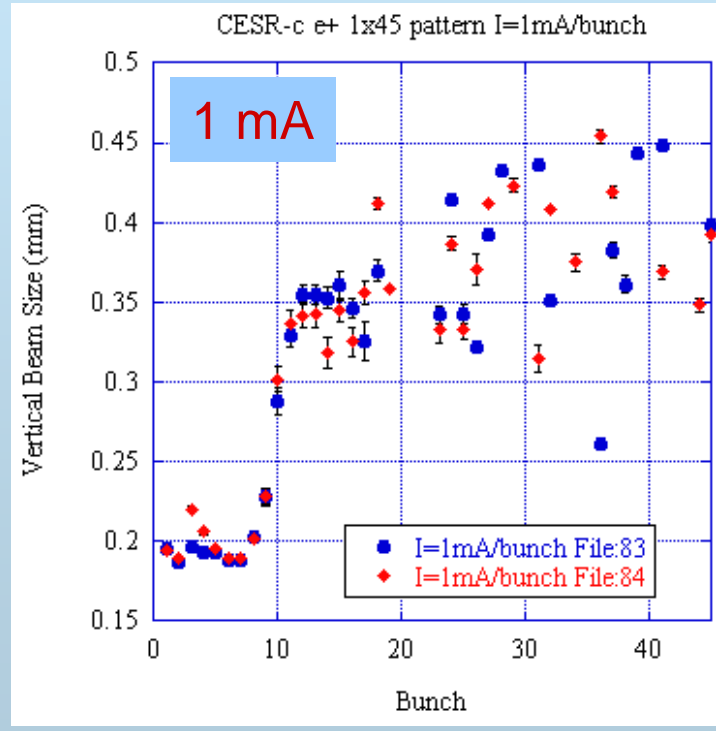
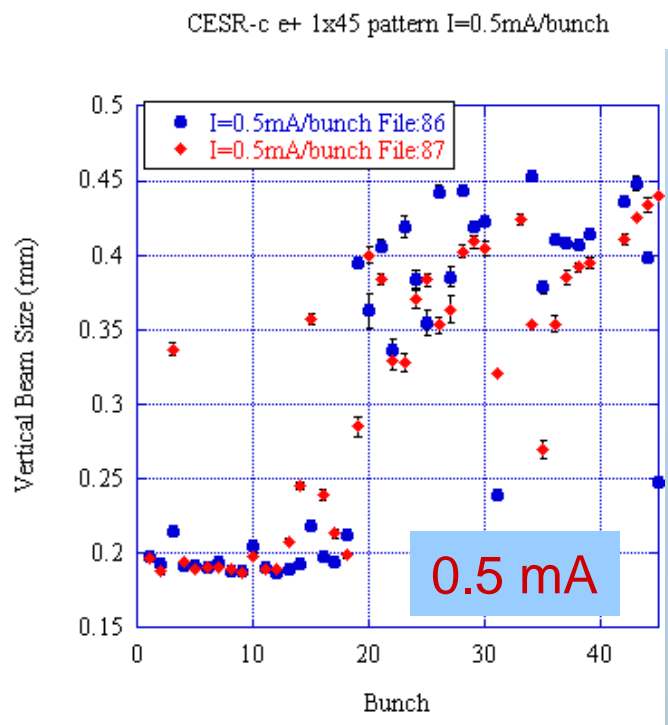
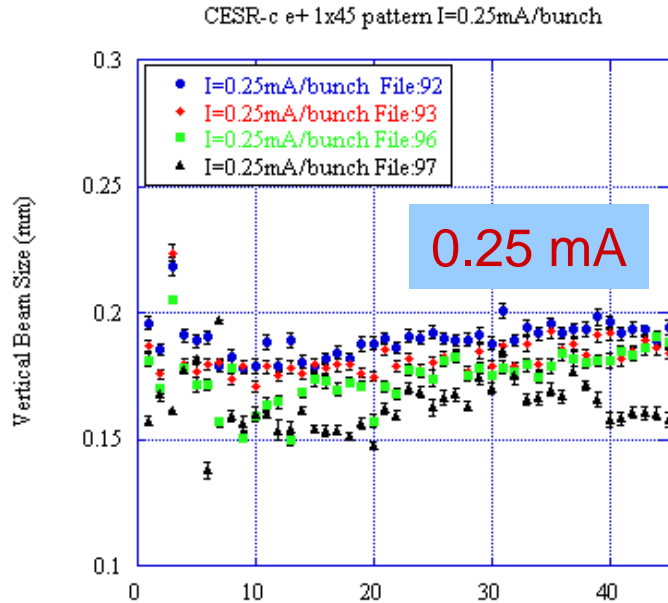


$$a = 4.3 \times 10^{-3}$$

$$T = 1.4 \times 10^{-8}$$

$$f_p = 1.483 \times 10^8$$

Vertical Beam-size for 1x45 Pattern





Outline

- *Photo-Electrons*
- *Dynamics of Electrons in e-Cloud*
- *e-Cloud Spatial Distribution*
- *Wake in the e-Cloud*
- *Experimental Implications*

Experimental Implications

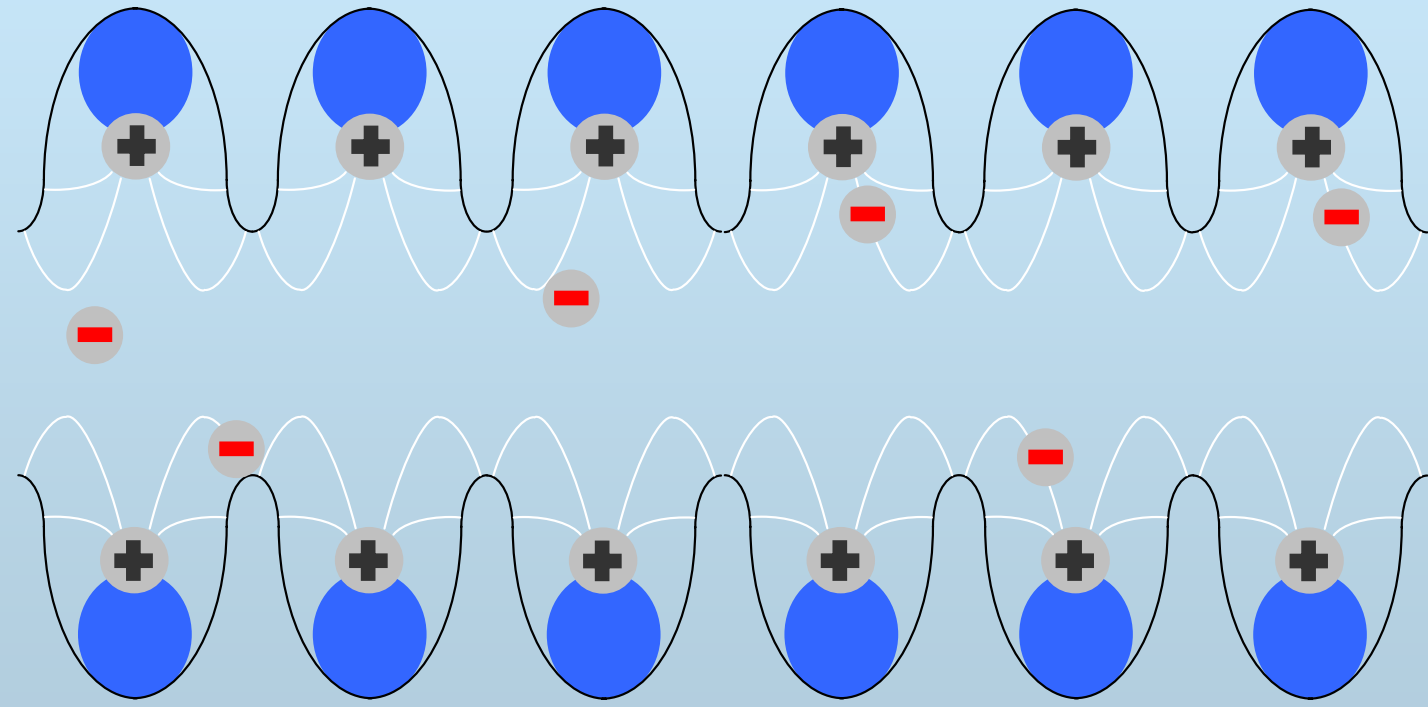
- Radiation generated in the *wiggler* is the main source of PE.
- *Life-time* of the electrons in the cloud plays crucial role.
- SE *extend* the life-time of the cloud.
- Relative tune shift *proportional* to the average cloud density.
- *Potential distribution*: main source of asymmetry between e^- & e^+ .
- “*Vacuum Cleaner*” may be efficient in particular for high N_{ec} .
Repels ions from the surface.
- *Space-charge* waves in the cloud generate *microwave* radiation.
- Radiation and transverse kick peak at *resonance* ($T=T_p$).
- SC wake treat differently trailing bunches.
- Wake has an impact also in case of an *array* of e-Clouds.



“Vacuum Cleaner”

Not in scale in order to emphasize the concept.

y ↑



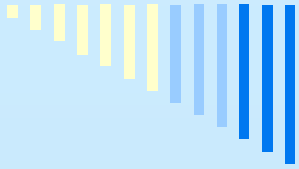


Photo-Electrons – Preliminary Conclusions

	$E_{th} [eV]$	$E_k = 2 [GeV]$		$E_k = 5 [GeV]$	
		$\rho_{eff} (E) = 0$	$\rho_{eff} (E) \neq 0$	$\rho_{eff} (E) = 0$	$\rho_{eff} (E) \neq 0$
N_{pe} / N_e	0	25.8	20.9	32.9	27.6
	300	0.047(0.18%)	0.047(0.22%)	1.1(3.3%)	1.1(4%)
V_{pe}	0	354.4[V]	332.9[V]	1.64[kV]	1.61[kV]
	300	19.8[V](5.6%)	19.8[V](6%)	0.97[kV](59%)	0.97[kV](60%)
N_{ph} / N_e	0	259.4		648.5	
	300	13.8(5%)		314 (48%)	
V_{ph}	0	177[kV]		0.690[MV]	
	300	6.6[kV] (4%)		0.667[MV](97%)	

- Number of PE is *one* order of magnitude smaller than the photons
- Effective voltage associated with PE *two* orders of magnitude smaller